

Twisted index theory on orbifold symmetric products and the fractional quantum Hall effect

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We extend the noncommutative geometry model of the fractional quantum Hall effect, previously developed by Mathai and the first author, to orbifold symmetric products. It retains the same properties of quantization of the Hall conductance at integer multiples of the fractional Satake orbifold Euler characteristics. We show that it also allows for interesting composite fermions and anyon representations, and possibly for Laughlin type wave functions.

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1. Introduction

A satisfactory model of the integer quantum Hall effect within the framework of noncommutative geometry was developed in [9], [10]. In a 2-dimensional periodic lattice, the presence of an external magnetic field turns the classical Brillouin zone into a non-commutative torus, replacing the ordinary translational symmetries of the Hamiltonian by magnetic translations. These are symmetries of the magnetic Laplacian, and they only commute up to a phase factor, hence the appearance of the noncommutative torus as the new algebra of observables. The integer quantization of the Hall conductance can then be interpreted in terms of an index theorem on the noncommutative torus.

In [49], [50], [51] a single particle model was developed for a charged particle moving in a magnetic field within a curved geometry described by a good 2-dimensional orbifold, with the curved geometry simulating an averaged effect of the interaction with other particles. This model exhibits quantization of the Hall conductance at fractional values given by integer multiples of the Satake orbifold Euler characteristic. The results of [49], [50], [51] are based on a generalization to the (fractional) orbifold case of a previous treatment, in [14], of the integer quantum Hall effect in the hyperbolic geometry of a smooth Riemann surfaces of genus $g \geq 2$.

The main drawback of this noncommutative-geometric approach to the fractional quantum Hall effect lies in the fact that it is still based on a single particle model. While the integer quantum Hall effect is well described by an independent electron approximation, which reduces it to a single particle model, the fractional quantum Hall effect is intrinsically a many particle phenomenon: while the single particle model used in [49], [50], [51] produces a fractional quantization of the Hall conductance as a Kawasaki orbifold index-theorem on the relevant noncommutative space, it does not account for Laughlin type wave functions, nor for composite fermion (or anyon) representations.

In this paper we propose a way to extend the noncommutative geometry model of [49], [50], [51], so that it makes contact with field theories on orbifolds, of the kind considered in relation to String Theory, see e.g. [1], [21], [54], [69]. We consider systems of n indistinguishable particles moving in the same type of homogeneous negatively curved geometry, under the effect of an external magnetic field, so that a classical configuration of the system is described by a point in the n -fold symmetric products of a good 2-dimensional orbifold. As in the case of field theories on orbifolds, the relevant Fock space is given by the sum of the orbifold K-theories (or the delocalized

equivariant cohomology) of the twisted group C^* -algebras of the wreath products $\Gamma_n = \Gamma^n \rtimes S_n$, with Γ the orbifold fundamental group of the good 2-dimensional orbifold. The quantization of the Hall conductance is still obtained via a twisted higher index theorem as in [50] and is expressible in terms of the Satake orbifold Euler characteristics of the orbifold symmetric products. At the same time, the model now allows for interesting composite fermion and anyon representations, whose classification depends on Seifert invariants of orbifold line bundles. We also formulate some hypothesis, still speculative at this stage, on how to obtain Laughlin type wave functions from the geometry of the model.

The paper is structured as follows: in the rest of this introductory section we discuss the geometry of 2-dimensional good orbifolds Σ and their symmetric products $\text{Sym}^n(\Sigma)$. In Section 2, we introduce the relevant groups that we will be considering, related to various kinds of orbifold covers. In particular, we extend to orbifold fundamental groups a result for smooth Riemann surfaces, which identifies the (orbifold) fundamental group of the symmetric products $\text{Sym}^n(\Sigma)$ with the abelianization of the (orbifold) fundamental group of Σ . In Section 3, we focus on K-theoretic aspects. We compute the orbifold K -theory of $\text{Sym}^n(\Sigma)$, in terms of classifying spaces for proper actions and we relate it to the K -theory of the group C^* -algebras $C_r^*(\Gamma_n)$ via the Baum–Connes conjecture, which we show is satisfied by the wreath products Γ_n . We also discuss a possible notion of orbifold-Jacobian and its K-theoretic properties. In Section 4, we show that the magnetic field determines a compatible family of $U(1)$ -multipliers σ_n on Γ_n . We obtain in this way twisted group C^* -algebras $C_r^*(\Gamma_n, \sigma_n)$ generalizing the algebra $C_r^*(\Gamma, \sigma)$ considered in [49], [50], [51]. Using the Lyndon–Hochschild–Serre spectral sequence for the group cohomology of the wreath products, we show that the multipliers σ_n define cocycles with trivial Dixmier–Douady class. We then show that the K -theory of the twisted group C^* -algebra $C_r^*(\Gamma_n, \sigma_n)$ agrees with the K -theory of the untwisted $C_r^*(\Gamma_n)$. To this purpose, we prove the K-amenability of the wreath products $\text{SL}(2, \mathbb{R})^n \rtimes S_n$, by adapting the argument of [30] for the K-amenability of $\text{SL}(2, \mathbb{R})$. Section 5 contains expository material, where we recall and compare the different notions of orbifold Euler characteristic used in the Kawasaki index theorem and in string theory on orbifolds. The former, which we refer to as the Satake orbifold Euler characteristic is in general a rational number, while the latter, which we call the string-theoretic orbifold Euler characteristic is an integer. We recall how the latter relates to sectors and to inertia orbifolds, and how it generalizes to orbifold Chern classes. We also recall the Segal construction of the Fock space for orbifold symmetric products, based on equivariant K-theory, as

in [64], [71]. In Section 6 we compute the twisted higher index theorem of [50] for the orbifold symmetric products $\text{Sym}^n(\Sigma)$. Using the same relation between Hall conductance cocycle and area cocycle as in [14], [50], we show that the Hall conductance is quantized at fractional values equal to integer multiples of the Satake orbifold Euler characteristic $\chi^{orb}(\text{Sym}^n(\Sigma))$. In Section 7 we classify composite fermions and anyons on the symmetric products $\text{Sym}^n(\Sigma)$. We introduce a notion of orbifold braid group, which is the orbifold fundamental group of the configuration spaces $\text{Conf}(\Sigma, n)$, where the orbifold singularities are coming from the orbifold cone points of Σ . We show that, similarly to what happens in the case of Riemann surfaces and ordinary braid groups, if the orbifold Σ has genus $g > 0$, then the scalar unitary representations can only be fermions or bosons, with no non-trivial anyons. When the genus is $g = 0$, there are anyon representations and we show that they are classified by the Seifert invariants of an orbifold line bundle with integer orbifold Euler number. We also show that, for arbitrary genus, there are anyon representations of higher dimensions N , which again depend on Seifert invariants, for an orbifold line bundle whose orbifold Euler number is in $\mathbb{Z} + (g + n - 1)/N$, and with fractional statistics $\pi i/N$. Finally Section 8 contains some more speculative considerations on how to find Laughlin type wave functions in this geometric setting, in terms of the Mathai–Quillen formalism for Euler classes of vector bundles, and local systems determined by a given N -dimensional anyon representation. The relation between N and $\#G$ imposed by the classification of anyon representations implies that one finds powers of the Vandermonde determinant with exponents equal to the denominators that appear in the quantization of the Hall conductance, as expected in Laughlin wave functions. We also suggest the possibility that Laughlin type functions may appear in computations via Selberg integrals of the orbifold Euler characteristic of a moduli space of good 2-dimensional orbifolds, analogous to the known calculations for moduli spaces of curves with marked points.

In the rest of this introductory section we review some known material that we need in the following.

1.1. Hyperbolic 2-dimensional good orbifolds

Let \mathbb{H} denote the 2-dimensional hyperbolic plane. We will use either the upper half plane model $\mathbb{H} \cong \{z = x + iy \in \mathbb{C} \mid \Im(z) > 0\}$ with the metric $ds^2 = (dx^2 + dy^2)/y^2$ or the equivalent Poincaré disc model $\mathbb{H} \cong \{z \in \mathbb{C} \mid |z| < 1\}$, with the metric $ds^2 = 4(dx^2 + dy^2)/(1 - |z|^2)^2$.

On \mathbb{H} we consider the isometric action of a discrete cocompact subgroup $\Gamma \subset \mathrm{PSL}(2, \mathbb{R})$, given by a Fuchsian group of signature $(g, \underline{\nu})$ with $\underline{\nu} = (\nu_1, \dots, \nu_m)$. These groups have an explicit presentation with generators a_i, b_i with $i = 1, \dots, g$ and c_j with $j = 1, \dots, m$, of the form

$$(1.1) \quad \Gamma = \Gamma(g, \underline{\nu}) = \left\langle a_i, b_i, c_j \left| \prod_{i=1}^g [a_i, b_i] c_1 \cdots c_m = 1, \ c_j^{\nu_j} = 1 \right. \right\rangle.$$

The quotient $\Sigma = \Sigma(g, \underline{\nu}) = \mathbb{H}/\Gamma$ is a 2-dimensional hyperbolic good orbifold. It is a Riemann surface of genus g with m cone points $\{x_1, \dots, x_m\}$, where the point x_j has stabilizer of order ν_j . Any such orbifold has a finite branched covering by a smooth Riemann surface $\Sigma_{g'}$, with $\Sigma(g, \underline{\nu}) = \Sigma_{g'}/G$ for a finite group G . The genus g' is related to g by the Riemann–Hurwitz formula for branched coverings:

$$(1.2) \quad g' = 1 + \frac{\#G}{2} \left(2(g-1) + \left(m - \sum_j \nu_j^{-1} \right) \right).$$

The Riemann surface $\Sigma_{g'} = \mathbb{H}/\Gamma'$ has a hyperbolic uniformization by $\Gamma' \subset \mathrm{PSL}(2, \mathbb{R})$, which is related to $\Gamma(g, \underline{\nu})$ by an exact sequence

$$(1.3) \quad 1 \longrightarrow \Gamma_{g'} \longrightarrow \Gamma(g, \underline{\nu}) \longrightarrow G \longrightarrow 1.$$

The Fuchsian group $\Gamma = \Gamma(g, \underline{\nu})$ is the *orbifold fundamental group* of the 2-dimensional orbifold $\Sigma = \Sigma(g, \underline{\nu})$, see [62].

1.2. Symmetric products

Let $\mathrm{Sym}^n(X) = X^n/S_n$, with S_n the group of permutations of a set of n elements. It is well known that, for a smooth compact Riemann surface Σ of genus g , the symmetric products $\mathrm{Sym}^n(\Sigma)$ are smooth and are related to the Jacobian of Σ in the following way. Let $\Omega^1(\Sigma)$ be the space of holomorphic 1-forms with a basis $\{\omega_1, \dots, \omega_g\}$. The group \mathcal{P}_Σ of periods of Σ is the subgroup $\mathcal{P}_\Sigma \subset \mathbb{C}^g$ given by the $v = (v_1, \dots, v_g) \in \mathbb{C}^g$ obtained as integrals $v_i = \int_\gamma \omega_i$ for some $\gamma \in \pi_1(\Sigma)$. For a base point $x_0 \in \Sigma$, and a path γ' from x_0 to $x \in \Sigma$, the integral $\int_{\gamma'} \omega_i$ then defines the Abel–Jacobi map

$$(1.4) \quad \mathcal{A} : \Sigma \rightarrow J(\Sigma) = \mathbb{C}^g/\mathcal{P}_\Sigma, \quad \mathcal{A} : x \mapsto \left(\int_{x_0}^x \omega_1, \dots, \int_{x_0}^x \omega_g \right),$$

where the Jacobian $J(\Sigma)$ is a torus T^{2g} , which can also be identified with $H^1(\Sigma, \mathbb{R})/H^1(\Sigma, \mathbb{Z})$. The Abel–Jacobi map extends to a map

$$(1.5) \quad \mathcal{A} : \text{Sym}^n(\Sigma) \rightarrow J(\Sigma), \quad \mathcal{A} : [x_1, \dots, x_n] \mapsto \mathcal{A}(x_1) + \dots + \mathcal{A}(x_n).$$

When $n > 2g - 2$, the symmetric products fiber over the Jacobian with fibers that are projective spaces of dimension $n - g$,

$$(1.6) \quad \mathbb{P}^{n-g}(\mathbb{C}) \hookrightarrow \text{Sym}^n(\Sigma) \twoheadrightarrow J(\Sigma).$$

Moreover, it was shown in Remark 5.8 of [35] (see also [41] for a generalization) that the fundamental group of the symmetric products of a smooth compact Riemann surface satisfies, for all $n \geq 2$,

$$(1.7) \quad \pi_1(\text{Sym}^n(\Sigma)) = \pi_1(\Sigma)^{ab} = H^1(\Sigma, \mathbb{Z}) = \pi_1(J(\Sigma)).$$

In the following section we will consider the symmetric products $\text{Sym}^n(\Sigma)$ of a 2-dimensional hyperbolic orbifold $\Sigma = \Sigma(g, \underline{\nu})$ and we compute the orbifold fundamental group.

2. Symmetric products, orbifold fundamental group, and orbifold coverings

In this section we discuss various orbifold coverings of the symmetric products $\text{Sym}^n(\Sigma)$ and their associated groups of symmetries.

Lemma 2.1. *Let Σ be a good 2-dimensional orbifold, with singular locus Σ_{sing} given by a finite set of cone points, and with orbifold fundamental group $\pi_1^{\text{orb}}(\Sigma)$. Then*

$$(2.1) \quad \pi_1^{\text{orb}}(\Sigma^n) \cong \pi_1^{\text{orb}}(\Sigma)^n.$$

Proof. Let \mathcal{O} be a good orbifold, with singular locus of (real) codimension two, $\text{codim} \mathcal{O}_{\text{sing}} = 2$. Then the orbifold fundamental group $\pi_1^{\text{orb}}(\mathcal{O})$ of an orbifold \mathcal{O} can be described ([68], §13) as the quotient

$$(2.2) \quad \pi_1^{\text{orb}}(\mathcal{O}) = \pi_1(\mathcal{O}_{\text{reg}})/H,$$

of the fundamental group $\pi_1(\mathcal{O}_{\text{reg}})$ of the regular part $\mathcal{O}_{\text{reg}} = \mathcal{O} \setminus \mathcal{O}_{\text{sing}}$ of the orbifold (the complement of the singular locus) by the normal subgroup H generated by the classes $\gamma_j^{\nu_j}$ in $\pi_1(\mathcal{O}_{\text{reg}})$, where γ_j are loops around a

component \mathcal{O}_j of $\mathcal{O}_{\text{sing}}$ with ν_j the order of the stabilizer G_j of \mathcal{O}_j . In particular, for a 2-dimensional good orbifold Σ we have $\Sigma_{\text{sing}} = \{x_j\}_{j=1,\dots,m}$ the cone points with stabilizers $\mathbb{Z}/\nu_j\mathbb{Z}$. In the product Σ^n we have $(\Sigma^n)_{\text{sing}} = \bigcup_{k=1}^n \Sigma_{\text{sing},(k)}$, where $\Sigma_{\text{sing},(k)}$ means a copy of Σ_{sing} in the k -th factor and the full space Σ in all the other factors $\Sigma_{\text{sing},(k)} = \Sigma \times \cdots \times \Sigma_{\text{sing}} \times \cdots \times \Sigma$. Thus, the set of regular points $(\Sigma^n)_{\text{reg}} = \Sigma^n \setminus \Sigma_{\text{sing}}^n$ is given by $(\Sigma^n)_{\text{reg}} = (\Sigma_{\text{reg}})^n$, and we can unambiguously use the notation Σ_{reg}^n for this locus. We have $\pi_1(\Sigma_{\text{reg}}^n) = \pi_1(\Sigma_{\text{reg}})^n$. The normal subgroup H of $\pi_1(\Sigma_{\text{reg}}^n)$ is generated by loops $\gamma_{j,k}$ in Σ^n that circle around the components $\Sigma \times \cdots \times \{x_j\} \times \cdots \times \Sigma$ of $\Sigma_{\text{sing},(k)}$. It suffices then to observe that for $k \neq k'$ and for all j, j' , the elements $\gamma_{j,k}$ and $\gamma_{j',k'}$ commute in $\pi_1(\Sigma_{\text{reg}}^n)$, so that the group H is a direct product $H = \prod_{k=1}^n H_k$ with each H_k isomorphic to the normal subgroup of the k -th factor $\pi_1(\Sigma_{\text{reg}})/H_k = \pi_1^{\text{orb}}(\Sigma)$. Thus, we obtain

$$\pi_1^{\text{orb}}(\Sigma^n) = \pi_1(\Sigma_{\text{reg}}^n)/H = \pi_1(\Sigma_{\text{reg}})^n / \prod_k H_k = \pi_1^{\text{orb}}(\Sigma)^n. \quad \square$$

Let Γ be a discrete group. Let $\Gamma^n \rtimes S_n$ be the semidirect product with multiplication

$$(g_1, \dots, g_n, \sigma)(h_1, \dots, h_n, \tau) = (g_1 h_{\sigma(1)}, \dots, g_n h_{\sigma(n)}, \sigma\tau),$$

and let $\langle S_n \rangle \subset \Gamma^n \rtimes S_n$ be the normal subgroup generated by the elements of S_n .

The following argument is implicit in Remark 5.8 of [35]. We reformulate it here in purely topological terms.

Lemma 2.2. *Let Σ be a good 2-dimensional orbifold, with singular locus Σ_{sing} given by a finite set of cone points. Then the orbifold fundamental groups of the symmetric products $\text{Sym}^n(\Sigma)$ satisfy*

$$(2.3) \quad \pi_1^{\text{orb}}(\text{Sym}^n(\Sigma)) \cong \pi_1^{\text{orb}}(\Sigma)^n \rtimes S_n / \langle S_n \rangle.$$

Proof. As in (2.2) in the previous Lemma, we have $\pi_1^{\text{orb}}(\Sigma) = \pi_1(\Sigma_{\text{reg}})/H$. The ordinary fundamental group $\pi_1(\Sigma_{\text{reg}})$ classifies (ordinary) covering spaces of Σ_{reg} , in the sense that, to each normal subgroup N of $\pi_1(\Sigma_{\text{reg}})$ there corresponds a regular covering space Σ_N of Σ_{reg} . Such a covering space is a principal $\pi_1(\Sigma_{\text{reg}})/N$ -fibration over Σ_{reg} . In particular, the quotient $\pi_1^{\text{orb}}(\Sigma) = \pi_1(\Sigma_{\text{reg}})/H$ similarly classifies all such coverings of Σ_{reg} that extend to a branched covering of Σ with a trivial action of the stabilizers of the singular points on the corresponding fibers. Such coverings correspond to

normal subgroups N of $\pi_1(\Sigma_{\text{reg}})$ that contain the normal subgroup H generated by loops γ_j around the components of Σ_{sing} with the appropriate multiplicities ν_j as above. Next observe that a regular covering Σ_N of Σ_{reg} can also be described as an S_n -equivariant covering $\tilde{\Sigma}_N$ of $\Sigma_{\text{reg}}^n \setminus \Delta$, the complement of the diagonals in $(\Sigma_{\text{reg}})^n = (\Sigma^n)_{\text{reg}}$. This means that there is an action of S_n on $\tilde{\Sigma}_N$, which is compatible with the action on $\Sigma_{\text{reg}}^n \setminus \Delta$, so that the projection map of the covering is S_n -equivariant. These are classified by the crossed product $\pi_1(\Sigma_{\text{reg}}^n \setminus \Delta) \rtimes S_n$ of the action of S_n on the fundamental group $\pi_1(\Sigma_{\text{reg}}^n \setminus \Delta)$. Among all such coverings, we consider those that extend to an S_n -equivariant branched covering of Σ^n , with a trivial action of the stabilizers of the components of Σ_{sing} and of the diagonals. These are then classified by a quotient of $\pi_1(\Sigma_{\text{reg}}^n \setminus \Delta) \rtimes S_n$ where we mod out by the normal subgroup generated by the loops around the components of Σ_{sing} and the elements of S_n , that is, by the group $\pi_1^{\text{orb}}(\Sigma^n) \rtimes S_n / \langle S_n \rangle$. Finally, observe that the data of an S_n -equivariant branched covering of Σ^n as above uniquely determine a branched covering over the symmetric product $\text{Sym}^n(\Sigma)$ and vice versa, so that we can identify $\pi_1^{\text{orb}}(\Sigma^n) \rtimes S_n / \langle S_n \rangle = \pi_1^{\text{orb}}(\text{Sym}^n(\Sigma))$. \square

The following observation is also implicit in Remark 5.8 of [35]. We spell it out for convenience.

Lemma 2.3. *Let $\Gamma^n \rtimes S_n$ be as above, with $\langle S_n \rangle \subset \Gamma^n \rtimes S_n$ the normal subgroup generated by the elements of S_n . There is a group isomorphism*

$$(2.4) \quad \Gamma^n \rtimes S_n / \langle S_n \rangle \cong \Gamma^{ab},$$

where $\Gamma^{ab} = \Gamma / [\Gamma, \Gamma]$ is the abelianization.

Proof. Let $g_{(i)}$ denote the element $g_{(i)} = (1, \dots, 1, g, 1, \dots, 1)$ of Γ^n that has the i -th coordinate equal to $g \in \Gamma$ and all the other coordinates equal to the identity element 1. Since Γ^n is a direct product of copies of Γ , the different factors Γ commute with each other, hence for all $g, h \in \Gamma$, we have $g_{(i)}h_{(j)} = h_{(j)}g_{(i)}$ whenever $i \neq j$, with the product equal to the element of Γ^n with g in the i -th place, h in the j -th place and 1 everywhere else. Moreover, observe that, for any $g \in \Gamma$ and for $i \neq j$, the element $(g_{(i)}, 1)$ in $\Gamma^n \rtimes S_n$ is equal to $(1, \sigma_{ij})^{-1}(g_{(j)}, 1)(1, \sigma_{ij})$, where $\sigma_{ij} \in S_n$ is the permutation that exchanges i and j and fixes all other elements of $\{1, \dots, n\}$. For any $\gamma \in \Gamma^n \rtimes S_n$, and for all $\sigma \in S_n$, we have $(1, \sigma)^{-1}\gamma(1, \sigma)\gamma^{-1} \in \langle S_n \rangle$, hence γ and $(1, \sigma)^{-1}\gamma(1, \sigma)$ define the same class in the quotient $\Gamma^n \rtimes S_n / \langle S_n \rangle$. In particular, $g_{(i)}$ and $g_{(j)}$ define the same element in the quotient, for all $g \in \Gamma$ and for all $i \neq j$. Thus, we obtain that in the quotient the n -copies of Γ in the product Γ^n

are identified and commutators are killed, hence the quotient gets identified with Γ^{ab} . \square

Combining the results of Lemmata 2.1, 2.2, and 2.3, we obtain the analog of (1.7) for orbifold fundamental groups.

Proposition 2.4. *For $n \geq 2$, the symmetric products $\text{Sym}^n(\Sigma)$ of a good 2-dimensional orbifold Σ have orbifold fundamental group given by*

$$(2.5) \quad \pi_1^{orb}(\text{Sym}^n(\Sigma)) \cong \pi_1^{orb}(\Sigma)^{ab}.$$

In the case of a 2-dimensional hyperbolic orbifold $\Sigma = \Sigma(g, \underline{\nu})$ one obtains the following.

Corollary 2.5. *For $\Sigma = \Sigma(g, \underline{\nu})$ and $n \geq 2$, we have $\pi_1^{orb}(\text{Sym}^n(\Sigma(g, \underline{\nu}))) = \mathbb{Z}^{2g} \oplus \mathbb{Z}_{\nu}$, where $\mathbb{Z}_{\nu} = \bigoplus_{j=1}^m \mathbb{Z}/\nu_j \mathbb{Z}$.*

Proof. The abelianization of a group $\Gamma = \Gamma(g, \underline{\nu})$ of the form (1.1) is given by $\mathbb{Z}^{2g} \oplus_j \mathbb{Z}/\nu_j \mathbb{Z}$. \square

Remark 2.6. The abelianization of the fundamental group is the first homology group, $\pi_1(\Sigma)^{ab} = H^1(\Sigma, \mathbb{Z})$. In the case of the orbifold fundamental group, there is a similar homological interpretation of its abelianization, in terms of the t -singular homology defined in [65], $\pi_1^{orb}(\Sigma)^{ab} = tH^1(\Sigma, \mathbb{Z})$, where the t -singular homology tH^* is constructed using singular simplexes that intersect transversely the singular locus of Σ , see [65] for details.

2.1. Geometry of some orbifold covering spaces

We consider some covering spaces of the orbifold symmetric products that will be useful in the rest of the paper.

Proposition 2.7. *For $n \geq 2$, let $\text{Sym}^n(\Sigma)$ be the symmetric product of the 2-dimensional orbifold $\Sigma = \Sigma(g, \underline{\nu})$. Let $\langle S_n \rangle$ be the normalizer of S_n in $\Gamma^n \rtimes S_n$ and let $\mathbb{S}_n = \langle S_n \rangle / S_n$. Let $\mathbb{S}^n(\mathbb{H}) := \mathbb{H}^n / \langle S_n \rangle$. Let $\Sigma_{g'} = \mathbb{H} / \Gamma_{g'}$ with $\Gamma_{g'}$ as in (1.3) with finite quotient $G = \Gamma / \Gamma_{g'}$. Let $G_n := G^n \rtimes S_n$.*

- 1) $\text{Sym}^n(\Sigma)$ is orbifold covered by \mathbb{H}^n , with $\text{Sym}^n(\Sigma) = \mathbb{H}^n / \Gamma^n \rtimes S_n$.
- 2) $\text{Sym}^n(\Sigma)$ is orbifold covered by $\mathbb{S}^n(\mathbb{H})$, with $\text{Sym}^n(\Sigma) = \mathbb{S}^n(\mathbb{H}) / \pi_1^{orb}(\text{Sym}^n(\Sigma))$.
- 3) $\mathbb{S}^n(\mathbb{H})$ is orbifold covered by $D^{2n} = \text{Sym}^n(\mathbb{H})$, with $\mathbb{S}^n(\mathbb{H}) = D^{2n} / S_n$.

- 4) $\text{Sym}^n(\Sigma)$ is orbifold covered by the smooth manifold $\Sigma_{g'}^n$, with $\text{Sym}^n(\Sigma) = \Sigma_{g'}^n/G_n$.

Proof. The orbifold $\Sigma = \Sigma(g, \nu)$ has a finite branched cover by a smooth surface $\Sigma_{g'}$, of genus (1.2), so that $\Sigma = \Sigma_{g'}/G$, with the finite group G as in (1.3). Moreover, Σ also has universal orbifold cover \mathbb{H} , with $\Sigma = \mathbb{H}/\Gamma$, for $\Gamma = \Gamma(g, \nu)$. The isometric action of Γ on \mathbb{H} induces an isometric action of $\Gamma^n \rtimes S_n$ on the n -fold product \mathbb{H}^n , with quotient $\mathbb{H}^n/\Gamma^n \rtimes S_n = \text{Sym}^n(\Sigma)$. Consider the normal subgroup $\langle S_n \rangle \subset \Gamma^n \rtimes S_n$. We can equivalently describe the quotient above as $(\mathbb{H}^n/\langle S_n \rangle)/(\Gamma^n \rtimes S_n/\langle S_n \rangle) = \text{Sym}^n(\Sigma)$. The group S_n is normal inside $\langle S_n \rangle$ with quotient \mathbb{S}_n , and we can further write the quotient $\mathbb{H}^n/\langle S_n \rangle = (\mathbb{H}^n/S_n)/\mathbb{S}_n = \text{Sym}^n(\mathbb{H})/\mathbb{S}_n$. We identify the hyperbolic plane \mathbb{H} with its Poincaré disc model $\mathbb{H} = D^2$ (the open unit disc in \mathbb{R}^2 with the hyperbolic metric). By Lemma 5 of [42], there is a homeomorphism of pairs

$$(D^{2n}, \partial D^{2n} = S^{2n-1}) \cong (\text{Sym}^n(D^2), \text{Sym}^n(\overline{D^2}) \setminus \text{Sym}^n(D^2)),$$

where D^{2n} is an open $2n$ -dimensional disc. Thus, we can identify $\text{Sym}^n(\mathbb{H}) \cong D^{2n}$, with the metric induced by the hyperbolic metric on \mathbb{H} , so that $\text{Sym}^n(\mathbb{H})/\mathbb{S}_n = D^{2n}/\mathbb{S}_n$. Finally, consider the sequence of groups (1.3). The normal embedding $\Gamma_{g'} \hookrightarrow \Gamma$ determines a normal embedding $\Gamma_{g'}^n \hookrightarrow \Gamma^n \rtimes S_n$. The quotient group can be identified with $G^n \rtimes S_n$, where $G = \Gamma/\Gamma_{g'}$. We then rewrite the quotient $\mathbb{H}^n/\Gamma^n \rtimes S_n = \text{Sym}^n(\Sigma)$ as $(\mathbb{H}^n/\Gamma_{g'}^n)/(\Gamma^n \rtimes S_n/\Gamma_{g'}^n) = \Sigma_{g'}^n/G_n$, with the finite group $G_n = G^n \rtimes S_n$. \square

3. Group algebras and K -theory

We now compute the orbifold K -theory groups of the symmetric products $\text{Sym}^n(\Sigma)$ and we discuss their relation to the K -theory of group C^* -algebras.

3.1. Orbifold K -theory

Let X be a good orbifold that is orbifold covered by a smooth manifold Y with $X = Y/G$. Then the orbifold K -theory of X is given by

$$(3.1) \quad K_{orb}^\bullet(X) = K_\bullet(C_0(Y) \rtimes G) = K_G^\bullet(Y),$$

the G -equivariant K -theory of Y .

We know from [26], [49] that, for a good 2-dimensional orbifold $\Sigma = \Sigma(g, \underline{\nu})$, with m cone points x_j with stabilizers of order ν_j , the orbifold K -theory is given by

$$(3.2) \quad K_{orb}^\bullet(\Sigma) = \begin{cases} \mathbb{Z}^{2-m+\nu} & \bullet = 0 \\ \mathbb{Z}^{2g} & \bullet = 1 \end{cases},$$

where $\nu = \sum_{j=1}^m \nu_j$.

For the symmetric products $\text{Sym}^n(\Sigma)$, using the covering (4) of Proposition 2.7, we obtain

$$(3.3) \quad K_{orb}^\bullet(\text{Sym}^n(\Sigma)) = K_\bullet((C(\Sigma_{g'}) \rtimes G)^{\otimes n} \rtimes S_n) = K_{G^n \rtimes S_n}^\bullet(\Sigma_{g'}^n),$$

where $K_\bullet(C(\Sigma_{g'}) \rtimes G) = K_{orb}^\bullet(\Sigma)$. Using the orbifold cover (1) of Proposition 2.7, we see that it can also be described as

$$(3.4) \quad K_{orb}^\bullet(\text{Sym}^n(\Sigma)) = K_{\Gamma^n \rtimes S_n}^\bullet(\mathbb{H}^n).$$

These descriptions of the orbifold K -theory of the symmetric products fall into a general framework for studying equivariant K -theory with respect to the action on powers X^n of the wreath products

$$(3.5) \quad G \sim S_n := G^n \rtimes S_n,$$

for a finite group G acting on a smooth manifold X . Several important properties of the equivariant K -theory groups $K_{G^n \rtimes S_n}^\bullet(X^n)$ were studied in [71]. We recall some of the main results of [71] and we apply them to our case, described as in (3.3).

3.2. Classifying space for proper action and assembly map

It is known from [6] that, to a locally compact group G one can associate a universal space for proper actions $\underline{E}G$, and a classifying space for proper actions given by the quotient $\underline{B}G = \underline{E}G/G$, so that there is a Kasparov assembly map from the equivariant K -homology groups with G -compact support $K_\bullet^G(\underline{E}G)$ to the K -theory of the reduced group C^* -algebra

$$(3.6) \quad \mu : K_j^G(\underline{E}G) \rightarrow K_j(C_r^*(G)),$$

which assigns to an abstract G -equivariant elliptic operator its index. The group G satisfies the Baum–Connes conjecture if the map (3.6) is an isomorphism.

The Baum–Connes conjecture (in fact the stronger form with coefficients) is implied by the Haagerup property, [53]. All finite groups satisfy the Haagerup property, and Fuchsian groups are also in the list of groups that are known to satisfy it, see [15], [53]. However, while it is known from [18] that the class of groups satisfying the Haagerup property is closed under wreath products, this only refers to “standard” wreath products $G \sim H := G^{(H)} \rtimes H$ where $G^{(H)} = \bigoplus_{h \in H} G$. A more general class of wreath products, which includes the case $G^n \rtimes S_n$ that we are interested in, is given by the “permutation wreath products” $G \sim_X H = G^{(X)} \rtimes H$, where X is an H -set and the action of H on $G^{(X)} = \bigoplus_{x \in X} G$ is by permuting indices $x \in X$ with the H action. As shown in [18], the Haagerup property is a lot more delicate for the case of the permutation wreath products.

However, for two groups G and H that both satisfy the Haagerup property, even if the more general permutation wreath products $G \sim_X H = G^{(X)} \rtimes H$ do not necessarily satisfy the Haagerup property, they all do satisfy the Baum–Connes conjecture. This follows from the general result of Oyono–Oyono on Baum–Connes for certain group extensions, [58]. Thus, we have the following property.

Lemma 3.1. *The groups $\Gamma^n \rtimes S_n$, with $\Gamma = \Gamma(g, \nu)$ a Fuchsian group, satisfy the Baum–Connes conjecture, hence the assembly map*

$$(3.7) \quad \mu : K_{orb}^\bullet(\mathrm{Sym}^n(\Sigma)) = K_{\Gamma^n \rtimes S_n}^\bullet(\mathbb{H}^n) \rightarrow K_\bullet(C_r^*(\Gamma^n \rtimes S_n))$$

is an isomorphism.

Proof. The general result of [58] implies that the wreath products $\Gamma^n \rtimes S_n$ satisfy the Baum–Connes conjecture. The result then follows by identifying $\mathbb{H}^n = \underline{E}(\Gamma^n \rtimes S_n)$ and $\mathrm{Sym}^n(\Sigma) = \underline{B}(\Gamma^n \rtimes S_n)$, with models for the universal and classifying space for proper actions, respectively. To see this, we can use the fact that if for a group G a G -space Y is a model of the universal space for proper actions $\underline{E}G$, and $G' \subset G$ is a subgroup, then Y is also a model of $\underline{E}G'$ (Corollary 1.9 of [6]) and that if G is a Lie group and K is the maximal compact subgroup, then a model of $\underline{E}G$ is given by the quotient G/K . A Lie group is virtually connected if it has only finitely many connected components. For any virtually connected Lie group the quotient G/K is diffeomorphic to a Euclidean space. We apply the above to the group $\mathrm{PSL}(2, \mathbb{R})^n \rtimes S_n$. \square

3.3. Orbifold K -theory and Lie group quotient

Consider as above the virtually connected Lie group $\mathrm{PSL}(2, \mathbb{R})^n \rtimes S_n$ and let \mathcal{K}_n denote its maximal compact subgroup, with quotient $\mathrm{PSL}(2, \mathbb{R})^n \rtimes S_n / \mathcal{K}_n \simeq \mathbb{H}^n$. The orbifold symmetric product is obtained as the double quotient

$$\mathrm{Sym}^n(\Sigma) = \Gamma_n \backslash \mathrm{PSL}(2, \mathbb{R})^n \rtimes S_n / \mathcal{K}_n,$$

where $\Gamma_n = \Gamma^n \rtimes S_n$. Let

$$(3.8) \quad \mathcal{P}_n := \Gamma_n \backslash \mathrm{PSL}(2, \mathbb{R})^n \rtimes S_n,$$

$$(3.9) \quad \hat{\mathcal{P}}_n := \Gamma_{g'}^n \backslash \mathrm{PSL}(2, \mathbb{R})^n \rtimes S_n.$$

Then we have the following.

Lemma 3.2. *Let \mathcal{P}_n , $\hat{\mathcal{P}}_n$, and \mathcal{K}_n be as above and let $G_n = G^n \rtimes S_n$, with $G = \Gamma / \Gamma_{g'}$. The algebras $C_0(\mathcal{P}_n) \rtimes \mathcal{K}_n$ and $C_0(\hat{\mathcal{P}}_n / \mathcal{K}_n) \rtimes G_n = C(\Sigma_{g'}^n) \rtimes G_n$ are strongly Morita equivalent.*

Proof. By (3.8) and (3.9), and the fact that $\Gamma_{g'}^n \subset \Gamma^n \rtimes S_n$ is a normal subgroup with quotient $G_n = G^n \rtimes S_n$, we obtain

$$\mathcal{P}_n = G_n \backslash \hat{\mathcal{P}}_n, \quad \text{and} \quad \hat{\mathcal{P}}_n / \mathcal{K}_n = \Gamma_{g'}^n \backslash \mathbb{H}^n = \Sigma_{g'}^n.$$

The Morita equivalence then follows as in Proposition 1.2 of [49], by applying [34]. \square

The orbifold C^* -algebra is defined in [26] as

$$C^*(\mathrm{Sym}^n(\Sigma)) = C(\mathcal{F}_n) \rtimes SO(2n),$$

where \mathcal{F}_n is the frame bundle of the orbifold tangent bundle of $\mathrm{Sym}^n(\Sigma)$. By the same argument of [26] it is shown to be strongly Morita equivalent to

$$C^*(\mathrm{Sym}^n(\Sigma)) \simeq C(\Sigma_{g'}^n) \rtimes G_n.$$

Thus, combining [26] with Proposition 2.7, we obtain the following.

Corollary 3.3. *The algebras $C^*(\mathrm{Sym}^n(\Sigma))$, $C(\Sigma_{g'}^n) \rtimes G_n$, $C_0(\mathbb{H}^n) \rtimes \Gamma_n$, $C_0(\mathbb{S}^n(\mathbb{H})) \rtimes \Gamma^{ab}$, and $C_0(\mathcal{P}_n) \rtimes \mathcal{K}_n$ are all strongly Morita equivalent.*

3.4. A notion of orbifold-Jacobian

Given a good 2-dimensional orbifold $\Sigma = \Sigma(g, \nu)$, we define the orbifold-Jacobian of Σ to be the product

$$(3.10) \quad J^{orb}(\Sigma) := J(\Sigma) \times \prod_{j=1}^m \mu_{\nu_j}$$

where $J(\Sigma) = H^1(\Sigma, \mathbb{R})/H^1(\Sigma, \mathbb{Z}) = \mathbb{T}^{2g}$, a real torus of rank $2g$, and μ_{ν_j} denotes the group of roots of unity of order ν_j . The group structure on $J^{orb}(\Sigma)$ is the direct product $\mathbb{T}^{2g} \times \prod_{j=1}^m \mu_{\nu_j}$.

For each cone point x_j , $j = 1, \dots, m$ on the orbifold Σ , let C_j be the boundary of a small disc in Σ centered at x_j . Let $x_{j,k}$ for $k = 1, \dots, \nu_j$ be points on C_j . For such a collection of base points, we define an orbifold-Abel-Jacobi map $\mathcal{A}^{orb} = \{\mathcal{A}_{j,k}\}$ with

$$\mathcal{A}_{j,k} : \Sigma \rightarrow \mathbb{T}^{2g} \times \{\zeta_{j,k}\}, \quad \mathcal{A}_{j,k} : \omega \mapsto \int_{x_{j,k}}^x \omega$$

where $\zeta_{j,k}$ are the roots of unity in μ_{ν_j} . This extends to an orbifold-Abel-Jacobi map $\mathcal{A}^{orb} : \text{Sym}^n(\Sigma) \rightarrow J^{orb}(\Sigma)$ by $\mathcal{A}^{orb}[x_1, \dots, x_n] = \mathcal{A}^{orb}(x_1) + \dots + \mathcal{A}^{orb}(x_n)$.

3.5. K-theory and the orbifold-Jacobian

The reduced group C^* -algebra $C_r^*(G)$ of a discrete group G is the norm closure in the algebra of bounded operators on $\ell^2(G)$ of the group ring $\mathbb{C}[G]$, acting via the left regular representation $L_g \xi(g') = \xi(g^{-1}g')$.

Lemma 3.4. *The reduced group C^* -algebra $C_r^*(\pi_1^{orb}(\text{Sym}^n(\Sigma)))$ has K -theory isomorphic to the topological K -theory of the orbifold-Jacobian $J^{orb}(\Sigma)$.*

Proof. The group $\pi_1^{orb}(\text{Sym}^n(\Sigma)) = \pi_1^{orb}(\Sigma)^{ab}$ is abelian. Thus, the K -theory of the reduced group C^* -algebra can be identified with the topological K -theory of the dual group, under Pontrjagin duality,

$$(3.11) \quad K_j(C_r^*(\pi_1^{orb}(\Sigma)^{ab})) \simeq K_j(\widehat{C(\pi_1^{orb}(\Sigma)^{ab})}).$$

We have $\pi_1^{orb}(\Sigma)^{ab} = \mathbb{Z}^{2g} \oplus \bigoplus_j \mathbb{Z}/\nu_j \mathbb{Z}$. The Pontrjagin dual of a direct sum of abelian groups is the direct product of the Pontrjagin duals. The dual of

\mathbb{Z}^{2g} is a $2g$ -dimensional real torus $\mathbb{T}^{2g} = S^1 \times \cdots \times S^1$, while for each finite group $\mathbb{Z}/\nu_j\mathbb{Z}$ the Pontrjagin dual is the subgroup $\mu_{\nu_j} \subset S^1$ of ν_j -th roots of unity, which can be identified again with $\mathbb{Z}/\nu_j\mathbb{Z}$. Thus, we obtain a direct product

$$\widehat{\pi_1^{orb}(\Sigma)^{ab}} = \mathbb{T}^{2g} \times \prod_j \mu_{\nu_j}. \quad \square$$

3.6. A homotopy theoretic version

Consider a smooth surface $\Sigma_{g'} = \mathbb{H}/\Gamma_{g'}$ uniformized by the hyperbolic plane \mathbb{H} . The surface $\Sigma_{g'}$ is a model of the classifying space for proper actions $\Sigma_{g'} = \underline{B}\Gamma_{g'}$, with $\mathbb{H} = \underline{E}\Gamma_{g'}$ a model for a universal space for proper actions, [6].

It is known (see Theorem 1.1 of [70]) that a group homomorphism $\alpha : \Gamma_1 \rightarrow \Gamma_2$ induces a commutative diagram

$$(3.12) \quad \begin{array}{ccc} K_{\bullet}^{\Gamma_1}(\underline{E}\Gamma_1) & \xrightarrow{\mu} & K_{\bullet}(C^*(\Gamma_1)) \\ \downarrow \alpha_* & & \downarrow \alpha_* \\ K_{\bullet}^{\Gamma_2}(\underline{E}\Gamma_2) & \xrightarrow{\mu} & K_{\bullet}(C^*(\Gamma_2)). \end{array}$$

The analogous statement for reduced algebras $C_r^*(\Gamma_i)$ holds in general only for monomorphisms (Corollary 1.2 of [70]), but since the geometric left-hand-side is always functorial, under the hypothesis that the Baum–Connes conjecture holds, then the right-hand-side would also be functorial for the reduced case, as observed in [70]. We focus on the case where $\Gamma_1 = \Gamma_{g'}$ and $\Gamma_2 = \Gamma_{g'}^{ab} = H^1(\Sigma_{g'}, \mathbb{Z})$, with $\alpha : \Gamma_{g'} \rightarrow \Gamma_{g'}^{ab}$ the quotient map. In this case, we know that the groups involved satisfy the Baum–Connes conjecture, and we can think of the left-hand-side of the diagram (3.12) as a kind of “homotopy-theoretic Abel–Jacobi map” from the K -homology of the curve $\Sigma_{g'}$ to that of its Jacobian,

$$(3.13) \quad \begin{array}{ccc} K_{\bullet}^{\Gamma_{g'}}(\mathbb{H}) \simeq K_{\bullet}(\Sigma_{g'}) & \xrightarrow{\mu} & K_{\bullet}(C^*(\Gamma_{g'})) \\ \downarrow \alpha_* & & \downarrow \alpha_* \\ K_{\bullet}^{\Gamma_{g'}^{ab}}(\mathbb{R}^{2g'}) \simeq K_{\bullet}(J(\Sigma_{g'})) & \xrightarrow{\mu} & K_{\bullet}(C^*(\mathbb{Z}^{2g'})). \end{array}$$

In a similar way, we obtain maps

$$(3.14) \quad K_{\bullet}^{\Gamma}(\mathbb{H}) \simeq K_{\bullet}^{orb}(\Sigma) \xrightarrow{\mu} K_{\bullet}(C^*(\Gamma)) \rightarrow K_{\bullet}(C^*(\Gamma^{ab})) = K^{\bullet}(J^{orb}(\Sigma))$$

and similar maps for the symmetric products

$$(3.15) \quad K_{\bullet}^{\Gamma^n \rtimes S_n}(\mathbb{H}^n) \simeq K_{\bullet}^{orb}(\mathrm{Sym}^n(\Sigma)) \xrightarrow{\mu} K_{\bullet}(C^*(\Gamma^n \rtimes S_n)) \\ \rightarrow K_{\bullet}(C^*(\Gamma^{ab})) = K^{\bullet}(J^{orb}(\Sigma)),$$

where the last map is induced by the quotient map

$$\Gamma^n \rtimes S_n \rightarrow \Gamma^n \rtimes S_n / \langle S_n \rangle \simeq \Gamma^{ab}.$$

4. The magnetic field and twisted group algebras

4.1. Twisted group ring and twisted group C^* -algebra

Recall that, for a discrete group Γ , a multiplier on Γ is defined as a map $\sigma : \Gamma \times \Gamma \rightarrow U(1)$ satisfying the properties:

- 1) $\sigma(\gamma, 1) = \sigma(1, \gamma) = 1$, for all $\gamma \in \Gamma$,
- 2) $\sigma(\gamma_1, \gamma_2)\sigma(\gamma_1\gamma_2, \gamma_3) = \sigma(\gamma_1, \gamma_2\gamma_3)\sigma(\gamma_2, \gamma_3)$, for all $\gamma_1, \gamma_2, \gamma_3 \in \Gamma$.

The reduced twisted group C^* -algebra $C_r^*(\Gamma, \sigma)$ is the norm closure in the algebra of bounded operators on $\ell^2(\Gamma)$ of the twisted group ring $\mathbb{C}[\Gamma, \sigma]$, generated by the left translations L_{γ}^{σ} with relations $L_{\gamma}^{\sigma}L_{\gamma'}^{\sigma} = \sigma(\gamma, \gamma')L_{\gamma\gamma'}^{\sigma}$, represented on $\ell^2(\Gamma)$ by the left regular representation

$$(4.1) \quad L_{\gamma}^{\sigma}\xi(\gamma') = \sigma(\gamma, \gamma^{-1}\gamma')\xi(\gamma^{-1}\gamma').$$

4.2. Cocycles from the magnetic field

In the system we are considering, we have n indistinguishable particles moving in a negatively curved geometry \mathbb{H} , subject to a potential, which is generated by charges disposed along the vertices of an embedding (Cayley graph) of a Fuchsian group Γ in \mathbb{H} , and to an external magnetic field.

Usually, in the “independent electron approximation” in the theory of solids, one replaces the (unbounded) interaction potential of the many-particle problem with a Hamiltonian with an effective potential of the form $H = \sum_{i=1}^n -\Delta_{x_i} + V(x_i)$, with a (bounded) effective potential $V(x)$ of a single-particle problem that encodes an average of the interactions of one of the charge carriers with the others, as well as with the atoms of the periodic medium. Here we consider a more general situation, where the independent electron approximation is perturbed by a (bounded) periodic interaction

potential. This means that we consider a Hamiltonian as above with the $V(x_i)$ replaced by a smooth bounded function $W(x_1, \dots, x_n)$ on \mathbb{H}^n , which is invariant under the symmetry group $\Gamma^n \rtimes S_n$.

As in [14], [49], [50], [51] the magnetic field is described by a closed 2-form $\omega = d\eta$ on \mathbb{H} . The form is invariant under the action of Γ on \mathbb{H} , so that $\omega - \gamma^*\omega = 0$ for all $\gamma \in \Gamma$. However, the potential η is not Γ -invariant. The fact that $d(\eta - \gamma^*\eta) = 0$ for all $\gamma \in \Gamma$ implies the existence of a (real valued) 0-form ϕ_γ on \mathbb{H} satisfying $\gamma^*\eta - \eta = d\phi_\gamma$, for all $\gamma \in \Gamma$. The function

$$(4.2) \quad \phi_\gamma = \int_{x_0}^x \gamma^*\eta - \eta$$

satisfies the property that

$$(4.3) \quad \phi_\gamma(x) + \phi_{\gamma'}(\gamma x) - \phi_{\gamma'\gamma}(x) = \phi_{\gamma'}(\gamma x_0)$$

is independent of $x \in \mathbb{H}$. Setting

$$(4.4) \quad \sigma(\gamma, \gamma') = \exp(-i\phi_{\gamma'}(\gamma x_0)),$$

for a chosen base point $x_0 \in \mathbb{H}$, determines a multiplier of Γ , as in §4.1. The magnetic Laplacian $\Delta^\eta = (d - i\eta)^*(d - i\eta)$ is invariant under the magnetic translations $T_\gamma^\phi = e^{-i\phi_\gamma} T_\gamma$ and the algebra of magnetic translations satisfies

$$(4.5) \quad T_\gamma^\phi T_{\gamma'}^\phi = \sigma(\gamma, \gamma') T_{\gamma\gamma'}^\phi,$$

with $\sigma(\gamma, \gamma')$ as in (4.4). This follows from (4.3).

Consider now the product \mathbb{H}^n and the 2-form $\varpi = \sum_j \omega_j$, where ω_j is the pullback $\omega_j = \pi_j^*\omega$ of the magnetic field 2-form described above, under the projection of \mathbb{H}^n onto the j -th factor. In particular ω_j only depends on the j -coordinate of \mathbb{H}^n .

Lemma 4.1. *The 2-form ϖ on \mathbb{H}^n is invariant under the action of $\Gamma_n = \Gamma^n \rtimes S_n$. The potential, given by the 1-form $\zeta = \sum_j \eta_j$ is invariant under S_n , hence it descends to a 1-form on $\text{Sym}^n(\mathbb{H}) = \mathbb{H}^n/S_n$. The form ζ , moreover, satisfies $g^*\zeta - \zeta = d\phi_g$, for $g = (\gamma, \sigma) \in \Gamma^n \rtimes S_n$, with*

$$(4.6) \quad \psi_g(x) = \sum_{j=1}^n \phi_{\gamma_j}(x_{\sigma(j)}),$$

with ϕ_{γ_j} as in (4.2) with a base point $x_{(0)} = (x_{0,j})$. This function $\psi_g : \mathbb{H}^n \rightarrow \mathbb{R}$ satisfies

$$(4.7) \quad \psi_g(x) + \psi_{g'}(gx) - \psi_{g'g}(x) = \psi_{g'}(gx_{(0)}),$$

independent of $x = (x_j) \in \mathbb{H}^n$. This determines a multiplier $\sigma_n : \Gamma_n \times \Gamma_n \rightarrow U(1)$,

$$(4.8) \quad \sigma_n(g, g') = \exp(-i\phi_{g'}(g x_{(0)})).$$

For $g \in \Gamma_n$, the magnetic translations T_g^ψ on \mathbb{H}^n satisfy

$$(4.9) \quad T_{g'}^\psi T_g^\psi = \sigma_n(g, g') T_{g'g}^\psi.$$

Proof. The identity (4.3) implies that for all $j = 1, \dots, n$ and all $\gamma_j \in \Gamma$, and all $\sigma, \sigma' \in S_n$ we have

$$\phi_{\gamma_{\sigma'(j)}}(x_{\sigma'\sigma(j)}) + \phi_{\gamma'_j}(\gamma_{\sigma'(j)} x_{\sigma'\sigma(j)}) - \phi_{\gamma'_j \gamma_{\sigma'(j)}}(x_{\sigma'\sigma(j)}) = \phi_{\gamma'_j}(\gamma_{\sigma'(j)} x_{0, \sigma'\sigma(j)}),$$

so that by summing over j we obtain (4.7). The composition of two magnetic translations then gives

$$e^{-i\psi_g(x)} e^{-i\psi_{g'}(gx)} f(g'gx) = e^{i(\psi_g(x) + \psi_{g'}(gx) - \psi_{g'g}(x))} T_{g'g}^\psi f(x),$$

so we obtain (4.9) with the multiplier (4.8). \square

4.3. K -theory of the twisted group algebra

The multiplier $\sigma_n : \Gamma_n \times \Gamma_n \rightarrow U(1)$ described above determines a twisting $C_r^*(\Gamma_n, \sigma_n)$ of the group C^* -algebra. The twisted group algebra provides the algebra of observables for the n -particle system, in the presence of the external magnetic field.

In [49] it was shown that the K -theory of the twisted group algebra $C_r^*(\Gamma, \sigma)$ is isomorphic to the K -theory of the untwisted algebra, whenever the multiplier $\sigma : \Gamma \times \Gamma \rightarrow U(1)$ has trivial Dixmier–Douady class $\delta(\sigma) = 0$. In the case of the algebras $C_r^*(\Gamma_n, \sigma_n)$, we have a similar results, as we will now discuss.

Proposition 4.2. *Let $\sigma : \Gamma \times \Gamma \rightarrow U(1)$ be a multiplier, with trivial Dixmier–Douady class, $\delta[\sigma] = 0$. Then it determines a multiplier $\sigma_n : \Gamma_n \times \Gamma_n \rightarrow U(1)$, which also has trivial Dixmier–Douady invariants, $\delta[\sigma_n] = 0$.*

Proof. Recall that the exponential sequence

$$1 \rightarrow \mathbb{Z} \xrightarrow{\iota} \mathbb{R} \xrightarrow{e} U(1) \rightarrow 1,$$

with $e(t) = \exp(2\pi it)$, determines a long exact cohomology sequence

$$\begin{aligned} \cdots \rightarrow H^2(\Gamma, \mathbb{Z}) \xrightarrow{\iota_*} H^2(\Gamma, \mathbb{R}) \xrightarrow{e_*} H^2(\Gamma, U(1)) \\ \xrightarrow{\delta} H^3(\Gamma, \mathbb{Z}) \xrightarrow{\iota_*} H^3(\Gamma, \mathbb{R}) \rightarrow \cdots \end{aligned}$$

where $\delta : H^2(\Gamma, U(1)) \rightarrow H^3(\Gamma, \mathbb{Z})$ is the Dixmier–Douady map. The Fuchsian group $\Gamma = \Gamma(g, \underline{\nu})$ has cohomology (see [49])

$$(4.10) \quad H^j(\Gamma, \mathbb{R}) = \begin{cases} \mathbb{R} & j = 0, 2 \\ \mathbb{R}^{2g} & j = 1 \\ 0 & j \geq 3 \end{cases}$$

A multiplier $\sigma : \Gamma \times \Gamma \rightarrow U(1)$ determines a cocycle $\sigma \in Z^2(\Gamma, U(1))$, with cohomology class $[\sigma] \in H^2(\Gamma, U(1))$. If the Dixmier–Douady class $\delta[\sigma] = 0$ in $H^3(\Gamma, \mathbb{Z})$, then the class $[\sigma]$ is in the range of the map $e_* : H^2(\Gamma, \mathbb{R}) \rightarrow H^2(\Gamma, U(1))$, that is, there exists a cocycle $\xi \in Z^2(\Gamma, \mathbb{R})$ such that $[\sigma] = [e(\xi)]$. Using the branched covering of the 2-dimensional orbifold Σ by a smooth Riemann surface $\Sigma_{g'}$ with $\Sigma = \Sigma_{g'}/G$, for a finite group G , we can identify $H^2(\Gamma, \mathbb{R}) \cong H^2(\Sigma_{g'}, \mathbb{R})$, since the finite group G has no nontrivial cohomology with real coefficients. Thus, we can realize the cocycle ξ in terms of a closed 2-form ω on $\Sigma_{g'}$, or of its $\Gamma_{g'}$ -invariant lift to the universal cover \mathbb{H} , with $[\sigma] = [e(\omega)]$. Consider then the cohomology $H^2(\Gamma_n, \mathbb{R})$. By the results of [47] [55] we know that the Lyndon–Hochschild–Serre spectral sequence for the group cohomology of the wreath product $\Gamma^n \rtimes S_n$ degenerates at the E_2 -term, for both integral cohomology and cohomology with coefficients in a field. In particular, this means that we can compute the cohomology of Γ_n with real coefficients in terms of the cohomology groups $E_2^{pq} = H^p(S_n, H^q(\Gamma^n, \mathbb{R}))$, with $H^q(\Gamma^n, \mathbb{R}) = \bigoplus_{i_1 + \cdots + i_n = q} H^{i_1}(\Gamma, \mathbb{R}) \times \cdots \times H^{i_n}(\Gamma, \mathbb{R})$. For $p + q = 2$, the only non-trivial term is $H^0(S_n, H^2(\Gamma^n, \mathbb{R}))$, since for the symmetric group $H^j(S_n, \mathbb{R}) = 0$ for $j \geq 1$. There is a subspace in $H^0(S_n, H^2(\Gamma^n, \mathbb{R})) \cong H^2(\Gamma^n, \mathbb{R})$ that is isomorphic to $H^2(\Gamma, \mathbb{R})^{\oplus n}$, namely the subspace given by the Künneth components involving only H^2 and H^0 and not H^1 . This subspace can be identified with $H^2(\Sigma_{g'}, \mathbb{R})^{\oplus n}$, as above. Using this identification, we see that the closed 2-form $\varpi(x) = \sum_j \omega(x_j)$ on $\Sigma_{g'}^n$ determines a class $[\varpi]$ in this subspace of $H^2(\Gamma_n, \mathbb{R})$, hence it determines

a multiplier class $[\sigma_n] = [e(\varpi)] \in H^2(\Gamma_n, U(1))$, as the image under the map e_* in the cohomology exact sequence

$$(4.11) \quad \begin{aligned} \cdots \rightarrow H^2(\Gamma_n, \mathbb{Z}) &\xrightarrow{\iota_*} H^2(\Gamma_n, \mathbb{R}) \xrightarrow{e_*} H^2(\Gamma_n, U(1)) \\ &\xrightarrow{\delta} H^3(\Gamma_n, \mathbb{Z}) \xrightarrow{\iota_*} H^3(\Gamma_n, \mathbb{R}) \rightarrow \cdots \end{aligned}$$

By construction, a representative $\sigma_n : \Gamma_n \times \Gamma_n \rightarrow U(1)$ of this class will be a multiplier with trivial Dixmier–Douady class, $\delta[\sigma_n] = 0$. \square

Remark 4.3. In the case of the multiplier $\sigma : \Gamma \times \Gamma \rightarrow U(1)$ defined by the magnetic field, we can take, in the argument of Lemma 4.2, the Γ -invariant 2-form ω on \mathbb{H} given by the magnetic field. The corresponding multiplier $\sigma_n : \Gamma_n \times \Gamma_n \rightarrow U(1)$ will then agree with the one constructed in the previous subsection. Thus, the σ and the σ_n determined by the magnetic field have trivial Dixmier–Douady invariant.

With the K -amenability property discussed in §4.4 below, we have the following.

Proposition 4.4. *Let $\sigma_n : \Gamma_n \times \Gamma_n \rightarrow U(1)$ be a multiplier as above, with $\delta[\sigma_n] = 0$. Then $K_\bullet(C_r^*(\Gamma_n, \sigma_n)) \simeq K_\bullet(C_r^*(\Gamma_n))$.*

Proof. The argument is the same as in [49]. If $\delta[\sigma_n] = 0$, we have $[\sigma_n] = [e(\xi_n)]$ for some $\xi_n \in Z^2(\Gamma_n, \mathbb{R})$, by (4.11). We can then use a homotopy $[\sigma_{n,t}] = [e(t\xi_n)]$ with $0 \leq t \leq 1$. Consider the discrete subgroup Γ_n of $\mathcal{G}_n = \mathrm{PSL}(2, \mathbb{R})^n \rtimes S_n$, with quotient $\Gamma_n \backslash \mathcal{G}_n = \mathcal{P}_n$ and let \mathcal{A} be an algebra with an action of Γ_n by automorphisms. The crossed product $(\mathcal{A} \otimes C_0(\mathcal{G}_n)) \rtimes \Gamma_n$ is Morita equivalent to the algebra of sections $C_0(\Gamma_n \backslash \mathcal{G}_n, \mathcal{E})$ of the flat \mathcal{A} -bundle $\mathcal{E} \rightarrow \hat{\mathcal{P}}_n$ with $\mathcal{E} = (\mathcal{A} \times \mathcal{G}_n)/\Gamma_n$ with the quotient taken with respect to the diagonal action. Moreover, the algebras $(\mathcal{A} \rtimes \Gamma_n) \otimes C_0(\mathcal{G}_n)$ and $(\mathcal{A} \otimes C_0(\mathcal{G}_n)) \rtimes \Gamma_n$ have the same \mathcal{K}_n -equivariant K -theory. Combined with the previous Morita equivalence and the fact that $\mathcal{G}_n/\mathcal{K}_n = \mathbb{H}^n$, we obtain

$$K_{\mathcal{K}_n, \bullet}(C_0(\mathcal{P}_n, \mathcal{E})) \cong K_{\mathcal{K}_n, \bullet + \dim(\mathcal{G}_n/\mathcal{K}_n)}(\mathcal{A} \rtimes \Gamma_n) = K_{\mathcal{K}_n, \bullet}(\mathcal{A} \rtimes \Gamma_n).$$

As in [49], we use the Packer–Raeburn stabilization trick [59]. The algebra $\mathcal{A} \rtimes_{\sigma_n} \Gamma_n$ is stably isomorphic to $(\mathcal{A} \otimes \mathbb{K}) \rtimes \Gamma$ with \mathbb{K} the algebra of compact

operators, and we consider the flat $\mathcal{A} \otimes \mathbb{K}$ -bundle

$$\mathcal{E}_{\sigma_n} = (\mathcal{A} \otimes \mathbb{K} \times \mathcal{G}_n) / \Gamma_n \rightarrow \Gamma_n \backslash \mathcal{G}_n.$$

As in Proposition 2.2 of [49], we then have

$$K_{\bullet}(C^*(\Gamma_n, \sigma_n)) \cong K_{\mathcal{K}}^{\bullet}(\Gamma_n \backslash \mathcal{G}_n, \delta(B_{\sigma_n})).$$

Here the twisted \mathcal{K}_n -equivariant K -theory $K_{\mathcal{K}}^{\bullet}(\Gamma_n \backslash \mathcal{G}_n, \delta(B_{\sigma_n}))$ is the same as the \mathcal{K}_n -equivariant K -theory of the continuous trace C^* -algebra $B_{\sigma_n} = C_0(\Gamma_n \backslash \mathcal{G}_n, \mathcal{E}_{\sigma_n})$ with Dixmier–Douady class $\delta(B_{\sigma_n})$. By Theorem 2.3 of [49], using the K -amenability property of §4.4 below, we then have $K_{\bullet}(C^*(\Gamma_n, \sigma_n)) \cong K_{\bullet}(C_r^*(\Gamma_n, \sigma_n))$. We then obtain isomorphisms $K_{\bullet}(C_r^*(\Gamma_n, \sigma_{n,t})) = K_{\bullet}(C_r^*(\Gamma_n))$. \square

Combining the isomorphism $K_{\bullet}(C_r^*(\Gamma_n, \sigma_{n,t})) = K_{\bullet}(C_r^*(\Gamma_n))$ obtained above with the Kasparov assembly map, we obtain a *twisted Kasparov map* (as in [49], [50])

$$(4.12) \quad \mu_{\sigma_n} : K_{orb}^{\bullet}(\text{Sym}^n(\Sigma)) \rightarrow K_{\bullet}(C_r^*(\Gamma_n, \sigma_n)).$$

4.4. K-amenability

We now turn to the K-amenability property of the group $\text{SL}(2, \mathbb{R})^n \rtimes S_n$. We first recall some basic facts about K-amenability. A locally compact second countable group \mathcal{G} is amenable if the map $\lambda : C^*(\mathcal{G}) \rightarrow C_r^*(\mathcal{G})$ determined by restriction of representations is an isomorphism. In particular, if \mathcal{G} acts on a C^* -algebra A as a C^* -dynamical system, and \mathcal{G} is amenable, then the map $\lambda_A : C^*(G, A) \rightarrow C_r^*(G, A)$ is also an isomorphism. The notion of K-amenability expresses a weaker K-theoretic form of this property. We consider here two forms of the K-amenability property, as in [30]. Recall that a Fredholm \mathcal{G} -module is a pair (\mathcal{H}, F) of a Hilbert space $\mathcal{H} = \mathcal{H}_0 \oplus \mathcal{H}_1$ with unitary representations ρ_0, ρ_1 of \mathcal{G} on \mathcal{H}_0 and \mathcal{H}_1 and with a bounded operator $F : \mathcal{H}_0 \rightarrow \mathcal{H}_1$ with

$$g \mapsto \rho_1(g) \circ F - F \circ \rho_0(g)$$

a compact operator and with $F^*F - 1$ and $FF^* - 1$ also compact operators. The trivial Fredholm module has $\mathcal{H}_0 = \mathbb{C}$ and $\mathcal{H}_1 = 0$. The notion of a homotopy of Fredholm module is similarly stated (see e.g. [30]).

- 1) \mathcal{G} is K-amenable if, for any C^* -dynamical system (\mathcal{G}, A) , the map $\lambda_{A*} : K_*(C^*(G, A)) \rightarrow K_*(C_r^*(G, A))$ is an isomorphism.
- 2) \mathcal{G} is K-amenable if there exists a Fredholm \mathcal{G} -module (\mathcal{H}, F) such that the representation of \mathcal{G} on the Hilbert space \mathcal{H} is weakly contained in the left regular representation, with (\mathcal{H}, F) homotopic to the trivial Fredholm \mathcal{G} -module.

The second version of the K-amenability property implies the first: we are going to refer to this second property as K-amenability. It was shown in [43] that any covering group of the identity component of $SO(n, 1)$ is K-amenable. In particular, $SL(2, \mathbb{R})$ is K-amenable. It is also shown in [19] that the class of K-amenable groups is closed with respect to the operations of taking subgroups and taking direct products. Thus, the groups $SL(2, \mathbb{R})^n$ are K-amenable.

Proposition 4.5. *The wreath product groups $SL(2, \mathbb{R})^n \rtimes S_n$ are K-amenable.*

Proof. The K-amenability of $SL(2, \mathbb{R})$ follows from the general result of [43] mentioned above. A more explicit proof was given in [30], by constructing a Fredholm module with the desired properties. This has $\mathcal{H}_0 = L^2(K/M)$ where K is the maximal compact (the circle group), with basis $\phi_n(\theta) = e^{in\theta}$, $n \in 2\mathbb{Z}$. and $M = \{\pm 1\} \subset K$, and $\mathcal{H}_1 = \mathcal{H}_{+2} \oplus \mathcal{H}_{-2}$ consisting of Hilbert space completions of the two discrete series representations, respectively given by the spans of $\{w_n \mid n \in 2\mathbb{Z}, n \geq 2\}$ and $\{w_n \mid n \in 2\mathbb{Z}, n \leq -2\}$, and with $F(\phi_0) = 0$ and $F(\phi_n) = w_n$. See §1 of [30] for more details. The K-amenability of $SL(2, \mathbb{R})^n$ follows, as mentioned above, from the general result of [19] which in particular shows the property is preserved by direct products. The construction of a Fredholm module for $SL(2, \mathbb{R})^n$ can be obtained from the construction of [30] by tensor products. Notice that the weak containment of representations has the property that, if unitary representations π_1 and π_2 are, respectively, weakly contained in unitary representations ρ_1 and ρ_2 , then $\pi_1 \otimes \pi_2$ is weakly contained in $\rho_1 \otimes \rho_2$. Let π be a unitary representation of a locally compact group \mathcal{G} on a Hilbert space \mathcal{H} . It induces a representation \mathcal{R}_π of the Banach algebra $L^1(\mathcal{G})$ on the same Hilbert space. Moreover, one has a \star -homomorphism $\mathcal{R} : L^1(\mathcal{G}) \rightarrow C_r^*(\mathcal{G})$. The representation π is weakly contained in the regular representation if $\|\mathcal{R}_\pi(f)\| \leq \|\mathcal{R}(f)\|$ for all $f \in L^1(\mathcal{G})$, see Definition 9.2.7 of [32]. For $\mathcal{G}_n = SL(2, \mathbb{R})^n \rtimes S_n$, consider the \mathcal{G}_n -Fredholm module given by $\mathcal{H}^{\otimes n}$, with \mathcal{H} the Fredholm module of [30]. The operator induced by F commutes with elements of S_n , and one

obtains in this way a \mathcal{G}_n -Fredholm module. The representation of \mathcal{G}_n is still weakly contained in the regular representation, because the inequality above is still satisfied for $f \in L^1(\mathcal{G}_n)$. \square

As an alternative, it may also be possible prove the result above by adapting the argument in Proposition 2.5 and Corollary 2.6 of [19].

5. Different notions of orbifold Euler characteristic

There are several different notions of orbifold Euler characteristic used in the literature. We will recall here some of the main versions and their relation. One of the main difference is that some orbifold Euler characteristics are rational valued, while other, even though they appear to be defined as fractions, are in fact integer valued. In particular, we are interested here in distinguishing between the Satake notion of (rational valued) orbifold Euler characteristic, which plays an important role in the noncommutative geometry approach to the fractional quantum Hall effect developed in [49], [50], [51], and the notion of orbifold Euler characteristic that arises naturally in string theory, [21], [69]. It was shown in [37], by a simple calculation, that the latter is integer valued. We restrict our attention here to the case of good orbifolds, which are global quotients, since the specific cases we intend to focus on, the symmetric products $\text{Sym}^n(\Sigma)$ of good 2-dimensional orbifolds, belong to this class: they are global quotients $\text{Sym}^n(\Sigma) = \Sigma_g^n / G^n \rtimes S_n$, as we discussed above.

Let X be a smooth manifold and G a finite group, acting on X with an orbifold quotient X/G . Then the Satake orbifold Euler characteristic, [60], which we simply write as $\chi^{orb}(X/G)$ is given by

$$(5.1) \quad \chi^{orb}(X/G) = \frac{1}{\#G} \chi(X) \in \mathbb{Q}.$$

In [60] a Gauss–Bonnet theorem is proved for orbifolds, where the usual topological Euler characteristic is replaced by the orbifold version (5.1), which is no longer, in general, an integer. The index theorems for elliptic operators on orbifolds proved in [44], [45] generalize the result of [60]. They were used in [49], [50] to obtain fractional values of the Hall conductance as values of a higher twisted index theorem modeled on [44] and on the higher index theorem of [17].

A different notion of orbifold Euler characteristic arises naturally in the context of string theory on orbifolds, [21], [69]. We will refer to it here as “string-theoretic orbifold Euler characteristic”, and we will denote it by

$\chi^{orb}(X, G)$, again assuming that the orbifold is a global quotient X/G of a smooth manifold by a finite group action. This version of the orbifold Euler characteristic is defined as

$$(5.2) \quad \chi^{orb}(X, G) = \frac{1}{\#G} \sum_{gh=hg} \chi(X^{\langle g, h \rangle}),$$

where the sum is over all pairs of commuting elements in G and $X^{\langle g, h \rangle}$ is the (common) fixed point set of g and h . Although from this definition this also appears to be rational valued, it is shown in [37] that the sum in (5.2) can be equivalently written as

$$(5.3) \quad \chi^{orb}(X, G) = \sum_{[g]} \chi(X^g / C(g)),$$

where now the sum is over conjugacy classes $[g]$ and $C(g)$ is the centralizer of g in G . The denominator $\#G$ disappears due to the simple fact that $\#[g] \cdot \#C(g) = \#G$. In orbifold string theory, the sum in (5.2) corresponds to the sum over the different sectors. Notice that the Satake orbifold Euler characteristic $\chi^{orb}(X/G)$ appears in the sum (5.2) as the term corresponding to the trivial sector with $g = h = 1$. It was shown in [4] that $\chi^{orb}(X, G) = \text{rank} K_G^0(X) - \text{rank} K_G^1(X)$, the difference of ranks of the equivariant K -theory.

Example 5.1. In the case of the 2-dimensional good orbifold $\Sigma = \Sigma_{g'}/G$ the Satake orbifold Euler characteristic is given by $\chi^{orb}(\Sigma) = (\#G)^{-1} \chi(\Sigma_{g'}) = \chi(\Sigma) + \sum_j (\nu_j^{-1} - 1)$, while the string-theoretic orbifold Euler characteristic is $\chi^{orb}(\Sigma_{g'}, G) = \chi(\Sigma_{g'}) + \sum_j (\nu_j - 1)$.

5.1. Orbifolds of A -sectors and inertia orbifolds

The Satake orbifold Euler characteristic and the string-theoretic orbifold Euler characteristic admit a family of common generalizations, see [66], [67] and in [27], [28], [29]. As above, let $Y = X/G$ be a good orbifold. We denote by $\mathcal{G} = \mathcal{G}(Y)$ the associated orbifold groupoid. In this setting one considers the additional data of a finitely generated discrete group A and defines the orbifold Y_A of A -sectors of Y through its orbifold groupoid, which is given by $\mathcal{G}(Y_A) = \mathcal{G}(Y) \ltimes \text{Hom}(A, \mathcal{G}(Y))$. One then defines the orbifold Euler characteristic $\chi_A^{orb}(X, G)$ as

$$(5.4) \quad \chi_A^{orb}(X, G) = \chi^{orb}(Y_A),$$

namely the Satake orbifold Euler characteristic of the orbifold Y_A . When $A = \mathbb{Z}^m$ one recovers the orbifold Euler characteristics $\chi_m^{orb}(X, G)$ of [13]. In particular, $\chi_{\mathbb{Z}^2}^{orb}(X, G) = \chi^{orb}(X, G)$, with $Y_{\mathbb{Z}^2}$ the inertia orbifold.

5.2. Orbifold Chern-Schwartz-MacPherson classes

A generalization of the generating function of (string theoretic) orbifold Euler characteristics (5.10) was given in [66], by considering the orbifold Euler characteristics $\chi_A^{orb}(X^n, G^n \rtimes S_n)$, with G a finite group. The case of the string-theoretic orbifold Euler characteristics of (5.2) and (5.10) is recovered for $A = \mathbb{Z}^2$. A further generalization of both (5.10) and the result of [66] was obtained in [56] as a generating function of orbifold characteristic classes, where the latter are defined as (equivariant) Chern-Schwartz-MacPherson classes, whose zero-dimensional component recovers the Euler characteristic. The notion of orbifold Chern-Schwartz-MacPherson classes considered in [56] is closely related to the *stringy Chern classes* of [2] and [20]. The orbifold CSM class is defined in [56] as the image under the equivariant MacPherson natural transformation $C_*^{G_n}$ (see [57]) of the canonical constructible function

$$\mathbf{1}_{\Sigma_{g'}^n, G_n}^{\mathbb{Z}^2} = \frac{1}{\#G_n} \sum_{\rho \in \text{Hom}(\mathbb{Z}^2, G_n)} \mathbf{1}_{(\Sigma_{g'}^n)^{\rho(\mathbb{Z}^2)}},$$

where $(\Sigma_{g'}^n)^{\rho(\mathbb{Z}^2)}$ is the fixed point set of the action of $\rho(\mathbb{Z}^2) \subset G_n$ on $\Sigma_{g'}^n$. Then the generating function of the orbifold CSM classes is then obtained by applying $C_*^{G, sym} = \prod_n C_*^{G_n}$ to the series $\sum_n \mathbf{1}_{\Sigma_{g'}^n, G_n}^{\mathbb{Z}^2} q^n$. This gives (Proposition 4.2 of [56])

$$\sum_n C_*^{G_n}(\Sigma_{g'}^n) q^n = \prod_{\ell=1}^{\infty} (1 - q^\ell \Delta^\ell)^{-C_*^G(\Sigma_{g'}^\ell)},$$

as in Theorem 1.2 of [56] with $A = \mathbb{Z}^2$, and with Δ^ℓ the morphism on homology induced by the diagonal embedding $\Delta : \Sigma_{g'} \hookrightarrow \Sigma_{g'}^\ell$. The string-theoretic orbifold Euler characteristics (5.10) are obtained by taking the 0-component of the CSM class, see (2.2) of [56]. Moreover, one can view the orbifold CSM class as in (2.2) and (2.4) of [56], as a sum

$$C_*^{orb}(X/G) = C_*(\pi_* \mathbf{1}_{X, G}^{\mathbb{Z}^2}) = \pi_* \iota^* C_*^G(\mathbf{1}_{X, G}^{\mathbb{Z}^2}) = \sum_g \iota_g^* C_*(X^g/C(g)),$$

where π maps G -invariant constructible functions on X to constructible functions on X/G and ι^* is the homomorphism from $H_*^G(X)$ to $H_*(X)$ (inclusion of X as fiber of $X \times_G EG \rightarrow BG$); the last sum is over conjugacy classes of elements $g \in G$, and $C(g)$ is the centralizer of g . In the latter form, these classes can be viewed (after reinterpreting them *cohomologically*) as residing in the *delocalized equivariant cohomology* $H^*(X, G) = \oplus H^*(X^g)^{C(g)}$, see [7]. Delocalized equivariant cohomology for symmetric products was considered in [72]. In the more general case of orbifolds X^n/G_n with a wreath product $G_n = G^n \rtimes S_n$, the delocalized equivariant cohomology is obtained as image under the Connes–Chern character of the K -theoretic construction of [71] recalled in §5.3 below. Delocalized equivariant cohomology is also the natural cohomology for string theory on orbifolds, in the sense of [21], [69].

Although the Chern–Schwartz–MacPherson classes are defined as homological Chern classes of *singular* varieties, they still admit a Chern–Weil type formulation in terms of curvature forms (currents), as shown in [31]. The construction of Chern–Weil representatives is based on an embedding of the singular variety X in a smooth ambient variety M , and in universal differential forms (currents) γ_k obtained from the pullbacks to $\mathbb{P}(T^*M)$ of the Chern classes $C_k(M)$ and the powers ζ^r of a 2-form ζ on $\mathbb{P}(T^*M)$ determined by the property that, on the total space $\mathbb{S}(T^*M)$ of the Hopf bundle $\pi_H : \mathbb{S}(T^*M) \rightarrow \mathbb{P}(T^*M)$ with fiber S^1 , one has $\pi_H^* \zeta = d\beta$, with β the generator of the cohomology of the fiber S^1 . More precisely, one has (§2 of [31])

$$\left(\sum_{r \geq 0} \zeta^r \right) \wedge \pi^* C_*(T^*M) = \sum_k (-1)^{\dim M - k} \gamma_k,$$

which, using the Chern–Weil curvature forms for $C_*(T^*M)$ leads to Chern–Weil representatives for the γ_k . The Chern–Schwartz–MacPherson class of $X \subset M$ is then obtained as the current $C_*(X)$ whose pairing with a form ω is given by

$$\langle C_k(X), \omega \rangle = \langle \mathbb{P}(N^*(X)), \gamma_k \wedge \omega \rangle = \int_{\mathbb{P}(N^*(X))} \gamma_k \wedge \omega,$$

where $N^*(X)$ is the Legendrian conormal cycle defined in §1 of [31] and $\mathbb{P}(N^*(X)) = \pi_{\mathcal{H}, \#}(N^*(X) \rfloor \beta)$, that is, $N^*(X)$ is the extension of $\mathbb{P}(N^*(X))$ to the total space of the Hopf bundle \mathcal{H} , $N^*(X) = \mathbb{P}(N^*(X)) \times_{\mathcal{H}} [S^1]$. In case of a smooth variety X one has $N^*(X) = (-1)^{\dim M - \dim X} [\mathbb{P}(\nu^*(X))]$ where $\nu^*(X)$ is the conormal bundle.

5.3. The Fock space of orbifold symmetric products

Following results of Segal for the equivariant K -theory of symmetric products [64], Weiqiang Wang established in [71] analogous results for the more general case of wreath products $G \sim S_n = G^n \rtimes S_n$, for a finite group G acting on a locally compact Hausdorff paracompact G -space X .

We use the notation $G_n = G^n \rtimes S_n$ as in [71]. Let $K_{G_n}^\bullet(X^n)$ be the equivariant K -theory and let $K_{G_n, \mathbb{C}}^\bullet(X^n) = K_{G_n}^\bullet(X^n) \otimes_{\mathbb{Z}} \mathbb{C}$. The Fock space $\mathcal{F}_G(X)$ is given by

$$(5.5) \quad \mathcal{F}_G(X) := \bigoplus_{n \geq 0} q^n K_{G^n \rtimes S_n, \mathbb{C}}^\bullet(X^n),$$

with q a formal variable (which keeps count of the graded structure) and with the term $n = 0$ equal to \mathbb{C} . It is proved in [71] that the Fock space $\mathcal{F}_G(X)$ has the following properties:

- $\mathcal{F}_G(X)$ is a graded connected Hopf algebra with multiplication defined by

$$(5.6) \quad K_{G_n, \mathbb{C}}^\bullet(X^n) \otimes K_{G_m, \mathbb{C}}^\bullet(X^m) \xrightarrow{K} K_{G_n \times G_m, \mathbb{C}}^\bullet(X^{n+m}) \xrightarrow{Ind} K_{G_{n+m}, \mathbb{C}}^\bullet(X^n),$$

where the first map K is the Künneth isomorphism and the second map Ind is the induction maps for subgroups; the comultiplication is given by

$$(5.7) \quad K_{G_n, \mathbb{C}}^\bullet(X^n) \xrightarrow{R} \bigoplus_{m=0}^n K_{G_m \times G_{n-m}, \mathbb{C}}^\bullet(X^n) \xrightarrow{K^{-1}} \bigoplus_{m=0}^n K_{G_m, \mathbb{C}}^\bullet(X^m) \otimes K_{G_{n-m}, \mathbb{C}}^\bullet(X^{n-m}),$$

where the first map R is restriction from G_n to subgroups $G_m \times G_{n-m}$ and the second map K^{-1} is the inverse of the Künneth isomorphism.

- As a graded algebra (graded over $\mathbb{Z}^+ \times \mathbb{Z}/2\mathbb{Z}$), $\mathcal{F}_G(X)$ is isomorphic to the algebra

$$(5.8) \quad \mathcal{S}(\bigoplus_{n \geq 1} q^n K_{G, \mathbb{C}}^0(X)) \otimes \Lambda(\bigoplus_{n \geq 1} q^n K_{G, \mathbb{C}}^1(X)),$$

where \mathcal{S} denotes the symmetric algebra and Λ the exterior algebra. The graded dimension satisfies

$$(5.9) \quad \dim_q \mathcal{F}_G(X) = \sum_{n \geq 0} q^n \dim K_{G_n, \mathbb{C}}(X^n) = \frac{\prod_{\ell \geq 1} (1 + q^\ell)^{\dim K_{G, \mathbb{C}}^1(X)}}{\prod_{\ell \geq 1} (1 - q^\ell)^{\dim K_{G, \mathbb{C}}^0(X)}}.$$

- The orbifold Euler characteristics (in the string theory sense) of the symmetric products have a generating function

$$(5.10) \quad \sum_{n \geq 1} \chi^{orb}(X^n, G_n) q^n = \prod_{\ell=1}^{\infty} (1 - q^\ell)^{-\chi^{orb}(X, G)}.$$

- $\mathcal{F}_G(X)$ is a free λ -ring generated by $K_{G, \mathbb{C}}^\bullet(X)$.

Remark 5.2. This construction, applied to $X = \Sigma_{g'}$ and $G = \Gamma(g, \underline{\nu})/\Gamma_{g'}$ provides a Fock space $\mathcal{F}_G(\Sigma_{g'})$ for our setting, with all the properties listed above, and

$$(5.11) \quad \sum_{n \geq 1} \chi^{orb}(\Sigma_{g'}^n, G^n \rtimes S_n) q^n = \prod_{\ell=1}^{\infty} (1 - q^\ell)^{-\chi^{orb}(\Sigma_{g'}, G)}.$$

6. Higher twisted index theory

Let \mathcal{E} be an orbifold vector bundle on the good 2-dimensional orbifold Σ . It defines a class $[\mathcal{E}]$ in $K_{orb}^\bullet(\Sigma)$. On the n -fold product Σ^n we consider the orbifold bundle $\mathcal{E}^{\boxtimes n}$. This determines an orbifold vector bundle \mathcal{E}_n on the symmetric product $\text{Sym}^n(\Sigma)$. Let $\tilde{\mathcal{E}}$ be the pull back of \mathcal{E} to \mathbb{H} and $\tilde{\mathcal{E}}^{\boxtimes n}$ the corresponding bundle on \mathbb{H}^n . Similarly, we consider the pullback \mathcal{E}' to $\Sigma_{g'}$ and the bundle $\mathcal{E}'^{\boxtimes n}$ on $\Sigma_{g'}^n$. The class $[\mathcal{E}_n] \in K_{orb}^\bullet(\text{Sym}^n(\Sigma))$ corresponds to the classes $[\tilde{\mathcal{E}}^{\boxtimes n}]$ and $[\mathcal{E}'^{\boxtimes n}]$, respectively, under the identifications $K_{orb}^\bullet(\text{Sym}^n(\Sigma)) = K_{\Gamma_n}^\bullet(\mathbb{H}^n) = K_{G_n}^\bullet(\Sigma_{g'}^n)$. Let $\partial_{\tilde{\mathcal{E}}}^+$ be the twisted Dirac operator on \mathbb{H} and let $\nabla = d - i\eta$ be the hermitian connection on \mathbb{H} with curvature $\nabla^2 = i\omega$, where ω is the Γ invariant 2-form defined by the magnetic field. The operator $\partial_{\tilde{\mathcal{E}}}^+ \otimes \nabla$ commutes with the projective action (Γ, σ) . Similarly, we write $\partial_{\mathcal{E}^{\boxtimes n}}^+$ for the twisted Dirac operator on \mathbb{H}^n and we consider $\partial_{\mathcal{E}^{\boxtimes n}}^+ \otimes \nabla_n$ where $\nabla_n = d - i\zeta$ where $\zeta = \sum_j \eta_j$ is the 1-form of Lemma 4.1. By the argument of Lemma 4.1, we see that $\partial_{\mathcal{E}^{\boxtimes n}}^+ \otimes \nabla_n$ commutes with the projective action (Γ_n, σ_n) . The analytic index is the image under the twisted Kasparov assembly map $\mu_{\sigma_n} : K_{orb}^\bullet(\text{Sym}^n(\Sigma)) \rightarrow K_\bullet(C_r^*(\Gamma_n, \sigma_n))$ of (4.12),

$$(6.1) \quad \text{Ind}_{(\Gamma_n, \sigma_n)}(\partial_{\mathcal{E}^{\boxtimes n}}^+ \otimes \nabla_n) = \mu_{\sigma_n}([\mathcal{E}_n]),$$

with the property that (§2.3 of [49])

$$(6.2) \quad \text{Ind}_{L^2}(\partial_{\mathcal{E}^{\boxtimes n}}^+ \otimes \nabla_n) = \text{tr}(\text{Ind}_{(\Gamma_n, \sigma_n)}(\partial_{\mathcal{E}^{\boxtimes n}}^+ \otimes \nabla_n)).$$

A cyclic 2-cocycle on an algebra \mathcal{R} is a multilinear map $t : \mathcal{R} \times \mathcal{R} \times \mathcal{R} \rightarrow \mathbb{C}$ satisfying

$$t(a, b, c) = t(c, a, b) = t(b, c, a)$$

$$t(ab, c, d) - t(a, bc, d) + t(a, b, cd) - t(da, b, c) = 0.$$

A dense involutive subalgebra $\mathcal{R}(\Gamma, \sigma) \subset C_r^*(\Gamma, \sigma)$, which contains the twisted group ring $\mathbb{C}[\Gamma, \sigma]$ and is closed under holomorphic functional calculus is constructed in §4 of [51], as the intersection of the domains of the powers δ^k of the derivation $\delta = [D, \cdot]$ associated to the operator $D\delta_\gamma = \ell(\gamma)\delta_\gamma$ that multiplies group elements by the word length $\ell(\gamma)$. The Haagerup inequality for surface groups shows that group cocycles on Γ with polynomial growth define cyclic cocycles on $\mathbb{C}[\Gamma, \sigma]$ that extend continuously to $\mathcal{R}(\Gamma, \sigma)$, see [51]. In particular, given a bounded 2-cocycle $c \in Z^2(\Gamma)$, one has an associated cyclic 2-cocycle tr_c on $\mathcal{R}(\Gamma, \sigma)$. This gives an additive map on K_0 , which we still denote by tr_c . Arguing as we did in §4 for the cocycle defined by the magnetic field, we can identify $H^2(\Gamma)^{\oplus n}$ with a subspace of $H^2(\Gamma_n)$, as in Proposition 4.2. Let c_n denote the cocycle in $H^2(\Gamma)^{\oplus n}$ defined by n copies of the 2-cocycle $c \in H^2(\Gamma)$, and let tr_{c_n} be the corresponding cyclic 2-cocycle on $\mathcal{R}(\Gamma_n, \sigma_n)$. As in [17] and in §3.2 of [50], one has an associated higher twisted analytic index

$$(6.3) \quad \begin{aligned} \text{Ind}_{(c_n, \Gamma_n, \sigma_n)}(\partial_{\mathcal{E}^{\boxtimes n}}^+ \otimes \nabla_n) &= \text{tr}_{c_n} \text{Ind}_{(\Gamma_n, \sigma_n)}(\partial_{\mathcal{E}^{\boxtimes n}}^+ \otimes \nabla_n) \\ &= \langle [\text{tr}_{c_n}], \mu_{\sigma_n}([\mathcal{E}_n]) \rangle = \langle [c_n], [\mathcal{E}_n] \rangle. \end{aligned}$$

We have the following index theorems (Theorem 1.1 of [49] and Theorem 2.2 of [50]), based on the Kawasaki index theorem on orbifold, [44], [45], see also [26], and on the higher index theorem of [17].

Proposition 6.1. *Suppose given a cocycle $c \in H^2(\Gamma)$ as above, with c_n the corresponding cocycle on Γ_n , together with a multiplier $\sigma : \Gamma \times \Gamma \rightarrow U(1)$ as in §4, determined by the closed Γ -invariant 2-form ω of the magnetic field, and the corresponding 2-form ω_n on $\Sigma_{g'}^n$. Then the higher twisted index (6.3) is given by*

$$(6.4) \quad \begin{aligned} &\text{tr}_{c_n} \text{Ind}_{(\Gamma_n, \sigma_n)}(\partial_{\mathcal{E}^{\boxtimes n}}^+ \otimes \nabla_n) \\ &= \frac{1}{(2\pi)^n n! (\#G)^n} \int_{\Sigma_{g'}^n} \hat{A}(\Omega_n) \text{tr}(e^{R^{\varepsilon' \boxtimes n}}) e^{\omega_n} \xi_{c_n}, \end{aligned}$$

where ξ_{c_n} is a 2-form representative of the class in $H^2(\Sigma_{g'}^n)^{\oplus n}$ corresponding to $c_n \in H^2(\Gamma)^{\oplus n} \subset H^2(\Gamma_n)$, and $n!(\#G)^n = \#G_n$.

The case (6.2), without the cyclic cocycle c , computes the range of the trace on K -theory, which is useful for gap labelling purposes, see [49], [51]. Here we focus on the higher version with the cyclic cocycle, as that will provide the quantization of the Hall conductance as in [51].

Lemma 6.2. *Let \mathcal{E} be an orbifold vector bundle over the good 2-dimensional orbifold $\Sigma = \Sigma_{g'}/G = \mathbb{H}/\Gamma$ and \mathcal{E}' the pull back to $\Sigma_{g'}$. Given a cocycle $c \in H^2(\Gamma)$ with the induced c_n on $\Gamma_n = \Gamma^n \rtimes S_n$, and let ω be the 2-form determined by the magnetic field. Then the twisted higher index theorem (6.4) can be written as*

$$(6.5) \quad \begin{aligned} & \mathrm{tr}_{c_n} \mathrm{Ind}_{(\Gamma_n, \sigma_n)} (\not{D}_{\mathcal{E}^{\boxtimes n}}^+ \otimes \nabla_n) \\ &= \frac{1}{(2\pi)^n n! (\#G)^n} \left(\int_{\Sigma_{g'}} \hat{A}(\Omega) \mathrm{tr}(e^{R^{\mathcal{E}'}}) e^{\omega} \xi_c \right)^n. \end{aligned}$$

Proof. The \hat{A} -genus is multiplicative over products, and the form $\hat{A}(\Omega_n)$ on $\Sigma_{g'}^n$ is the product of n copies of $\hat{A}(\Omega)$ on $\Sigma_{g'}$. The Chern character $\mathrm{ch}(\mathcal{E}) = \mathrm{tr}(e^{R^{\mathcal{E}}})$ is multiplicative for external tensor products, hence $\mathrm{tr}(e^{R^{\mathcal{E}' \boxtimes n}})$ on $\Sigma_{g'}^n$ is also a product of n copies of $\mathrm{tr}(e^{R^{\mathcal{E}'}})$ each depending only on one of the factors $\Sigma_{g'}$. The form ω_n is by construction (see §4) a sum of n copies of the 2-form ω , each depending only on the coordinates of one of the $\Sigma_{g'}$ factors, hence e^{ω_n} is also a product. Moreover, as we have seen above, the 2-form ξ_{c_n} on $\Sigma_{g'}^n$ is also a product of copies of a 2-form ξ_c on $\Sigma_{g'}$. Thus, the integrand in (6.4) splits as a product of identical terms depending on only one of the factors. \square

The area 2-cocycle $c \in H^2(\Gamma)$ is the restriction to $\Gamma \subset \mathrm{PSL}(2, \mathbb{R})$ of the hyperbolic area $c : \mathrm{PSL}(2, \mathbb{R}) \times \mathrm{PSL}(2, \mathbb{R}) \rightarrow \mathbb{R}$, where $c(\gamma_1, \gamma_2)$ is the oriented hyperbolic area of the geodesic triangle in \mathbb{H} with vertices $(z_0, \gamma_1^{-1}z_0, \gamma_2 z_0)$, for a chosen base point z_0 .

Corollary 6.3. *In the case where $c \in H^2(\Gamma)$ is the area 2-cocycle, the range of values of the twisted higher index theorem, while varying the choice of the orbifold vector bundle \mathcal{E} on Σ is given by*

$$(6.6) \quad \chi^{\mathrm{orb}}(\mathrm{Sym}^n(\Sigma)) \mathbb{Z} \subset \mathbb{Q},$$

where $\chi^{\mathrm{orb}}(\mathrm{Sym}^n(\Sigma))$ is the Satake orbifold Euler characteristic of $\mathrm{Sym}^n(\Sigma)$.

Proof. The integral

$$\frac{1}{2\pi\#G} \int_{\Sigma_{g'}} \hat{A}(\Omega) \text{tr}(e^{R^{\mathcal{E}'}}) e^{\omega} \xi_c$$

is the twisted higher index theorem computed in [50]. It is shown in Corollary 3.2 of [50] that, when $c \in H^2(\Gamma)$ is the area 2-cocycle, this integral is given by

$$\frac{\chi(\Sigma_{g'})}{\#G} \cdot \text{rank}(\mathcal{E}) = \chi^{orb}(\Sigma) \cdot \text{rank}(\mathcal{E}).$$

Thus, for the area cocycle, the twisted higher index theorem (6.5) is given by

$$\frac{\chi(\Sigma_{g'})^n}{n!(\#G)^n} \cdot \text{rank}(\mathcal{E})^n = \chi^{orb}(\text{Sym}^n(\Sigma)) \cdot \text{rank}(\mathcal{E})^n.$$

□

The Hall conductance on Σ is also described by a cyclic 2-cocycle on the twisted group algebra $\mathbb{C}[\Gamma, \sigma]$, given by

$$(6.7) \quad \text{tr}_K(f_0, f_1, f_2) = \sum_{j=1}^g \text{tr}(f_0(\delta_j(f_1)\delta_{j+g}(f_2) - \delta_{j+g}(f_1)\delta_j(f_2))),$$

where g is the genus and the δ_j , for $j = 1, \dots, 2g$ are derivations associated to the elements of a symplectic basis of $H^1(\Sigma, \mathbb{R})$. If P_E denotes the spectral projection associated to the Fermi level, then the Hall conductance is given by

$$\sigma_E = \text{tr}_K(P_E, P_E, P_E).$$

A derivation of this expression for the Hall conductance can be obtained as a quantum adiabatic limit, see [14]. As shown in Theorem 4.1 of [50], the conductance cocycle and the area cocycle are cohomologous. Since the twisted higher index theorem, seen as a pairing of cyclic cohomology and K-theory, only depends on the class of the cyclic cocycle, the range of the twisted higher index theorem also determines the possible range of values of the Hall conductance. We summarize the conclusion of this section as follows: the single particle theory on Σ with the external magnetic field ω extends in a compatible way to a many particles model on the symmetric products $\text{Sym}^n(\Sigma)$. In this model, the range of quantized values of the Hall conductance consists of integer multiples of the Satake orbifold Euler characteristics $\chi^{orb}(\text{Sym}^n(\Sigma))$.

7. Orbifold braid groups and anyons

In this section we analyze what types of anyons and composite fermions one obtains within this model of fractional quantum Hall effect. These are related to a notion of “orbifold braid groups” that we introduce below.

The configuration space of n (ordered) points on \mathbb{H} is given by the complement of the diagonals $F(\mathbb{H}, n) = \mathbb{H}^n \setminus \Delta$. The configuration spaces of unordered points is defined as the quotient by the action of the symmetric group

$$(7.1) \quad \text{Conf}(\mathbb{H}, n) := (\mathbb{H}^n \setminus \Delta)/S_n = F(\mathbb{H}, n)/S_n.$$

These have fundamental group $\pi_1(\text{Conf}(\mathbb{H}, n)) = B_n$, the Artin braid group, with generators σ_i , $i = 1, \dots, n-1$ and relations $\sigma_i \sigma_j = \sigma_j \sigma_i$ for $|i-j| \geq 2$ and $\sigma_i \sigma_{i+1} \sigma_i = \sigma_{i+1} \sigma_i \sigma_{i+1}$, for $i = 1, \dots, n-2$. In fact, the spaces $\text{Conf}(\mathbb{H}, n)$ are topologically Eilenberg-MacLane spaces $K(B_n, 1)$, see [23].

Given a 2-dimensional compact (topological) surface Σ , and a finite set of points $Q = \{x_j\}_{j=1, \dots, m}$ on Σ , one similarly defines the configuration spaces

$$(7.2) \quad \begin{aligned} F(\Sigma \setminus Q, n) &= (\Sigma \setminus Q)^n \setminus \Delta \\ \text{Conf}(\Sigma \setminus Q, n) &= F(\Sigma \setminus Q, n)/S_n. \end{aligned}$$

For $r < n$, the projections $\Pi_{n,r}(z_1, \dots, z_n) = (z_1, \dots, z_r)$ define locally trivial fibrations $F(\Sigma \setminus Q, n) \rightarrow F(\Sigma \setminus Q, r)$ with the fiber over $w = (w_1, \dots, w_r)$ given by the configuration space $F(\Sigma \setminus (Q \cup \{w_i\}_{i=1, \dots, r}), n-r)$. The braid group of $\Sigma \setminus Q$ on n strings is given by the fundamental group

$$(7.3) \quad B_n(\Sigma \setminus Q) := \pi_1(\text{Conf}(\Sigma \setminus Q, n)).$$

In particular, if Σ is a 2-dimensional orbifold and $Q \subset \Sigma$ is the set of cone points, we have corresponding braid groups

$$(7.4) \quad B_n(\Sigma_{\text{reg}}) = \pi_1(\text{Conf}(\Sigma_{\text{reg}}, n)) = \pi_1((\Sigma_{\text{reg}}^n \setminus \Delta)/S_n).$$

7.1. Orbifold braid group

For a good 2-dimensional orbifold Σ , we can also associate to the configuration space $\text{Conf}(\Sigma_{\text{reg}}, n)$ an *orbifold braid group*, defined as

$$(7.5) \quad B_n^{\text{orb}}(\Sigma) := \pi_1^{\text{orb}}(\text{Conf}(\Sigma, n)),$$

using the orbifold fundamental group of $\text{Conf}(\Sigma, n)$.

Proposition 7.1. *The orbifold braid group of Σ is a quotient of the ordinary braid group of Σ_{reg} by a normal subgroup generated by powers $\gamma_j^{\nu_j}$, with ν_j the order of the stabilizer of the j -th cone point x_j of Σ , and with γ_j a loop in $\text{Conf}(\Sigma_{\text{reg}}, n)$ that winds around the j -th component of the (real) codimension two stratum of $\text{Conf}(\Sigma, n)_{\text{sing}}$.*

Proof. We have $\text{Conf}(\Sigma_{\text{reg}}, n) = (\Sigma_{\text{reg}}^n \setminus \Delta)/S_n \subset \text{Sym}^n(\Sigma)_{\text{reg}}$, as in Lemma 2.2, where

$$\begin{aligned} (\Sigma^n \setminus \Delta)_{\text{reg}} &= (\Sigma^n \setminus \Delta) \cap \Sigma_{\text{reg}}^n \\ &= \{(z_1, \dots, z_n) \in \Sigma_{\text{reg}}^n \mid z_i \neq z_j \ \forall i, j\} \\ &= \{(z_1, \dots, z_n) \in \Sigma^n \mid z_i \in \Sigma_{\text{reg}} \text{ and } x_i \neq x_j \ \forall i, j\} \\ &= \Sigma_{\text{reg}}^n \setminus \Delta. \end{aligned}$$

Moreover, $((\Sigma_{\text{reg}}^n \setminus \Delta)/S_n)_{\text{reg}} = (\Sigma_{\text{reg}}^n \setminus \Delta)/S_n$, hence we can write the orbifold fundamental group as

$$\pi_1^{\text{orb}}(\text{Conf}(\Sigma, n)) = \pi_1(\text{Conf}(\Sigma_{\text{reg}}, n))/H,$$

where H is the normal subgroup generated by elements $\gamma_a^{\nu_a}$, where the γ_a are loops in $\text{Conf}(\Sigma_{\text{reg}}, n)$ around a component X_a of the singular locus $\text{Conf}(\Sigma, n)_{\text{sing}}$ with ν_a the order of the stabilizer of X_a . The singular locus $\text{Conf}(\Sigma, n)_{\text{sing}} = (\Sigma^n \setminus \Delta)_{\text{sing}}/S_n$ only comes from the cone points of Σ , namely

$$(\Sigma^n \setminus \Delta)_{\text{sing}} = (\Sigma^n)_{\text{sing}} \setminus \Delta,$$

where $(\Sigma^n)_{\text{sing}} = \cup_{k=1}^n \Sigma_{\text{sing}, k}$, with

$$\Sigma_{\text{sing}, k} = \Sigma \times \dots \times \Sigma \times \Sigma_{\text{sing}} \times \Sigma \dots \times \Sigma,$$

with a copy of Σ_{sing} in the k -th factor and Σ in all the other factors. We denote by $\Sigma_{\text{sing}, k}(x_j) \subset \Sigma_{\text{sing}, k}$ the component of the (real) codimension two stratum of $(\Sigma^n)_{\text{sing}}$ that has a cone point $\{x_j\} \subset \Sigma_{\text{sing}}$ in the k -th factor. Thus, the components of the (real) codimension two stratum of $(\Sigma^n)_{\text{sing}} \setminus \Delta$ are of the form $\Sigma_{\text{sing}, k}(x_j) \setminus \Delta$. Let $\gamma_{j, k}$ be a loop in $\Sigma^n \setminus \Delta$ that winds around the component $\Sigma_{\text{sing}, k}(x_j)$. The power $\gamma_{j, k}^{\nu_j}$, where ν_j is the order of the stabilizer of the cone point x_j in Σ is a generator of the subgroup H of

$\pi_1(\Sigma^n \setminus \Delta)$ such that

$$\pi_1^{orb}(\Sigma^n \setminus \Delta) = \pi_1(\Sigma^n \setminus \Delta)/H.$$

Passing to the quotient $F(\Sigma, n) = \Sigma^n \setminus \Delta \rightarrow \text{Conf}(\Sigma, n) = (\Sigma^n \setminus \Delta)/S_n$, the (real) codimension two stratum of the singular locus is the image of the S_n invariant configuration of components in $F(\Sigma, n)$ given by $\Sigma_{\text{sing}}^n(x_j) := \cup_{k=1}^n \Sigma_{\text{sing},k}(x_j)$, for a given cone point x_j . We denote by γ_j a loop in $\text{Conf}(\Sigma, n)$ winding around $\Sigma_{\text{sing}}^n(x_j)$. The powers $\gamma_j^{\nu_j}$ generate the subgroup H with

$$\pi_1^{orb}(\text{Conf}(\Sigma, n)) = \pi_1(\text{Conf}(\Sigma_{\text{reg}}, n))/H.$$

□

The braid groups of a 2-dimensional orientable (topological) surface of genus g with m punctures were computed explicitly in [11], [12] (see also [63], [8] for a slightly different form of the presentation). With the presentation given in [8], if Σ has genus g with a set $Q = \{x_j\}_{j=1,\dots,m}$ of m punctures, the braid group $B_n(\Sigma \setminus Q)$ has additional generators with respect to the Artin braid group B_n . Namely, the generators are given by

$$(7.6) \quad \begin{aligned} \sigma_i, & \quad i = 1, \dots, n-1 \\ a_\ell, & \quad \ell = 1, \dots, g \\ b_\ell, & \quad \ell = 1, \dots, g \\ c_j, & \quad j = 1, \dots, m-1 \end{aligned}$$

with relations

$$(7.7) \quad \begin{aligned} \sigma_i \sigma_j &= \sigma_j \sigma_i, & \text{when } |i-j| \geq 2 \\ \sigma_i \sigma_{i+1} \sigma_i &= \sigma_{i+1} \sigma_i \sigma_{i+1}, & i = 1, \dots, n-2, \end{aligned}$$

and in the Artin braid group, and additional relations, for all ℓ ,

$$(7.8) \quad \begin{aligned} a_\ell \sigma_i &= \sigma_i a_\ell, & i \neq 1, \\ b_\ell \sigma_i &= \sigma_i b_\ell, & i \neq 1, \end{aligned}$$

$$(7.9) \quad \begin{aligned} \sigma_1^{-1} a_\ell \sigma_1^{-1} a_\ell &= a_\ell \sigma_1^{-1} a_\ell \sigma_1^{-1} \\ \sigma_1^{-1} b_\ell \sigma_1^{-1} b_\ell &= b_\ell \sigma_1^{-1} b_\ell \sigma_1^{-1} \\ \sigma_1^{-1} a_\ell \sigma_1^{-1} b_\ell &= b_\ell \sigma_1^{-1} a_\ell \sigma_1 \\ \sigma_1^{-1} a_\ell \sigma_1 a_r &= a_r \sigma_1^{-1} a_\ell \sigma_1 & \ell < r, \\ \sigma_1^{-1} b_\ell \sigma_1 b_r &= b_r \sigma_1^{-1} b_\ell \sigma_1 & \ell < r, \\ \sigma_1^{-1} a_\ell \sigma_1 b_r &= b_r \sigma_1^{-1} a_\ell \sigma_1 & \ell < r, \\ \sigma_1^{-1} b_\ell \sigma_1 a_r &= a_r \sigma_1^{-1} b_\ell \sigma_1 & \ell < r, \end{aligned}$$

and for all ℓ, r and all j

$$(7.10) \quad \begin{aligned} c_j \sigma_i &= \sigma_i c_j, & i \neq 1, \\ \sigma_1^{-1} c_j \sigma_1 a_r &= a_r \sigma_1^{-1} c_j \sigma_1 & n > 1, \\ \sigma_1^{-1} c_j \sigma_1 b_r &= b_r \sigma_1^{-1} c_j \sigma_1 & n > 1, \\ \sigma_1^{-1} c_j \sigma_1 c_k &= c_k \sigma_1^{-1} c_j \sigma_1 & j < k, \\ \sigma_1^{-1} c_j \sigma_1^{-1} c_j &= c_j \sigma_1^{-1} c_j \sigma_1^{-1}. \end{aligned}$$

In the case without punctures, the braid group $B_n(\Sigma)$ has generators σ_i , a_ℓ , b_ℓ as above with the same relations (7.7), (7.8), (7.9), but with (7.10) replaced by

$$(7.11) \quad \prod_{\ell=1}^g [a_\ell, b_\ell^{-1}] = \sigma_1 \sigma_2 \cdots \sigma_{n-1}^2 \cdots \sigma_2 \sigma_1.$$

Proposition 7.2. *Let Σ be a 2-dimensional good orbifold of genus g , with m cone points. The orbifold braid groups $B_n^{\text{orb}}(\Sigma)$ has generators σ_i , $i = 1, \dots, n-1$, a_ℓ , b_ℓ , $\ell = 1, \dots, g$ and c_j , $j = 1, \dots, m$ with relations as above, and with the additional relation*

$$(7.12) \quad \prod_{\ell=1}^g [a_\ell, b_\ell^{-1}] \sigma_1^{-1} \sigma_2^{-1} \cdots \sigma_{n-1}^{-2} \cdots \sigma_2^{-1} \sigma_1^{-1} c_1 \cdots c_m = 1$$

and $c_j^{\nu_j} = 1$, where ν_j is the order of the stabilizer of the cone point.

Proof. As shown in [8], the generators a_ℓ and b_ℓ correspond geometrically, in terms of a fundamental domain for Σ given by a $4g$ -gons with pairwise identified sides marked by the generators α_ℓ and β_ℓ of $\pi_1(\Sigma)$. The braid a_ℓ is a string that crosses the α_ℓ sides and b_ℓ the β_ℓ side (with the opposite orientation), while the c_j wind around the j -th puncture. The generators σ_i have the usual meaning as in the Artin braid group. The relations are explained geometrically in §2.2 of [8], where it is also shown that one can equivalently introduce an additional generator c_m , as the braid that winds around the last puncture, and the additional relation (7.12). This corresponds to writing the fundamental group of the punctured surface $\Sigma \setminus Q = \Sigma_{\text{reg}}$ as

$$\pi_1(\Sigma \setminus Q) = \left\langle \{a_\ell, b_\ell\}_{\ell=1, \dots, g}, \{c_j\}_{j=1, \dots, m} \mid \prod_{\ell} [a_\ell, b_\ell^{-1}] c_1 \cdots c_m = 1 \right\rangle,$$

instead of writing it solely in terms of the generators a_ℓ, b_ℓ, c_j with $j = 1, \dots, m-1$. In particular, the generators c_j provide loops in the configuration space $\text{Conf}(\Sigma_{\text{reg}}, n)$ that wind around the component $\Sigma_{\text{sing}}^n(x_j)$. Thus,

by Proposition 7.1 we obtain $B_n^{orb}(\Sigma)$ from $B_n(\Sigma_{\text{reg}})$ by imposing the further relations $c_j^{\nu_j} = 1$. \square

7.2. Anyons

Fractional-statistics particles, or anyons, have the property that, when two particles get interchanged, the wavefunction changes by a phase factor $\exp(i\pi\alpha)$, for some $\alpha \in (-1, 1]$. The cases $\alpha = 0$ and $\alpha = 1$ correspond, respectively, to bosons and fermions. It is known [22], [38] (see also [48]) that, for 2-dimensional systems, the surface topology plays an important role in determining what type of anyons can arise. More precisely, these are classified by 1-dimensional unitary representations of the braid group of the surface. The case of the orbifold braid group is similar.

Lemma 7.3. *Let Σ be a 2-dimensional good orbifold of genus $g > 0$, with m cone points. One-dimensional unitary representations R of the orbifold braid group $B_n(\Sigma)$ have the generators σ_i acting as $R(\sigma_i) = \pm 1$, the generators a_ℓ and b_ℓ , respectively, acting as phase factors $R(a_\ell) = e^{2\pi i\theta_\ell}$ and $R(b_\ell) = e^{2\pi i\phi_\ell}$ and the generators c_j acting as $R(c_j) = e^{2\pi i\beta_j/\nu_j}$, where ν_j is the order of the stabilizer of the j -th cone point and the β_j satisfy $\sum_{j=1}^m \beta_j/\nu_j \in \mathbb{Z}$.*

Proof. The fact that the σ_i must act like ± 1 follows from the relation $\sigma_1^{-1}a_\ell\sigma_1^{-1}b_\ell = b_\ell\sigma_1^{-1}a_\ell\sigma_1$ in (7.9), which implies that σ_1 acts as ± 1 and the relations (7.7), which imply that all the σ_i must then also act in the same way. The action of the a_ℓ and b_ℓ is unconstrained by the relations, hence we get independent phase factors for each of them, while the c_j are constrained by the relations $c_j^{\nu_j} = 1$ and (7.12), which implies $R(c_1 \cdots c_m) = 1$. These give $R(c_j) = e^{2\pi i\beta_j/\nu_j}$ with $\sum_{j=1}^m \beta_j/\nu_j \in \mathbb{Z}$. \square

Remark 7.4. Since the braids σ_i correspond to exchanging two particles, all the representations in Lemma 7.3 are either fermions or bosons, whenever $g > 0$. The β_j can be viewed as the Seifert invariants of an orbifold line bundle over Σ that has integer orbifold Euler number (hence is an actual line bundle). See §7.3 for more details.

In the case of good 2-dimensional orbifolds of genus $g = 0$ with m cone points, we have the following result.

Lemma 7.5. *Let Σ be a 2-dimensional good orbifold of genus $g = 0$, with m cone points. Then the one-dimensional unitary representations R of the*

orbifold braid group $B_n(\Sigma)$ have $R(\sigma_i) = e^{i\pi\alpha}$ and $R(c_j) = e^{2\pi i\beta_j/\nu_j}$, where ν_j is the order of the stabilizer of the j -th cone point and the β_j satisfy

$$(7.13) \quad \alpha - \sum_{j=1}^m \frac{\beta_j}{\nu_j} \in \mathbb{Z}.$$

Proof. The argument is exactly as in the previous lemma, except that we do not have the generators a_ℓ , b_ℓ and the only relations are the (7.7) and

$$\begin{aligned} c_j \sigma_i &= \sigma_i c_j, & i &\neq 1 \\ \sigma_1^{-1} c_j \sigma_1 c_k &= c_k \sigma_1^{-1} c_j \sigma_1 & j < k, \\ \sigma_1^{-1} c_j \sigma_1^{-1} c_j &= c_j \sigma_1^{-1} c_j \sigma_1^{-1} \\ \sigma_1 \sigma_2 \cdots \sigma_{n-1}^2 \cdots \sigma_2 \sigma_1 &= c_1 \cdots c_m \end{aligned}$$

and $c_j^{\nu_j} = 1$. This last relation gives, as before, $R(c_j) = e^{2\pi i \frac{\beta_j}{\nu_j}}$. The relations (7.7) imply that all the σ_i must act by the same phase factor $R(\sigma_i) = e^{i\pi\alpha}$ and the last displayed relation then implies that $e^{2\pi i\alpha} = e^{2\pi i \sum_{j=1}^m \beta_j/\nu_j}$, hence we obtain (7.13). \square

Remark 7.6. Thus, in the case of good orbifolds of genus zero, there are non-trivial anyons (that are neither fermions nor bosons) and the fractional statistics they satisfy depends on the datum of an orbifold line bundle on Σ , through the Seifert invariants β_j , see §7.3 below.

Moreover, it is known that two-dimensional systems on surfaces of genus $g > 0$ do admit fractional statistics arising from higher dimensional irreducible unitary representations of the braid group $B_n(\Sigma)$, [22], [38], provided α satisfies $\exp(2\pi i(n + g - 1)\alpha) = 1$. We describe the analog for the orbifold braid group.

Consider the $N \times N$ matrices

$$(7.14) \quad U_N = \begin{pmatrix} 1 & & & & \\ & \xi_N^2 & & & \\ & & \xi_N^4 & & \\ & & & \ddots & \\ & & & & \xi_N^{2(N-1)} \end{pmatrix}, \quad \text{with} \quad \xi_N = \exp(\pi i/N),$$

$$(7.15) \quad V_N = \begin{pmatrix} 0 & 1 & 0 & \cdots & 0 & 0 \\ 0 & 0 & 1 & \cdots & 0 & 0 \\ \vdots & & & & & \vdots \\ 0 & 0 & 0 & \cdots & 0 & 1 \\ 1 & 0 & 0 & \cdots & 0 & 0 \end{pmatrix}.$$

They satisfy the commutation relation

$$(7.16) \quad V_N U_N = \xi_N^2 U_N V_N.$$

Proposition 7.7. *Let Σ be a good 2-dimensional orbifold of genus g with m cone points. The orbifold braid group $B_n^{\text{orb}}(\Sigma)$ has unitary representations of dimension N^g with $R(\sigma_i) = \xi_N^{-1}$ and $R(a_\ell) = U_{N,\ell}$ and $R(b_\ell) = V_{N,\ell}$, where $U_{N,\ell}$ and $V_{N,\ell}$ act as U_N and V_N , respectively, in the ℓ -th factor of the tensor product of g -copies of \mathbb{C}^N and the identity on the other factors. The generators c_j act as $R(c_j) = e^{2\pi i \beta_j / \nu_j}$, where ν_j is the order of the stabilizer of the j -th cone point and the β_j satisfy the relation*

$$(7.17) \quad \frac{(g+n-1)}{N} + \sum_{j=1}^m \frac{\beta_j}{\nu_j} \in \mathbb{Z}.$$

Proof. The relation $\sigma_1^{-1} a_\ell \sigma_1^{-1} b_\ell = b_\ell \sigma_1^{-1} a_\ell \sigma_1$ is now satisfied, since the left-hand-side $\xi_N^2 U_{N,\ell} V_{N,\ell}$ and the right-hand-side $V_N U_N$ agree by (7.16). The rest of the relations (7.9) are also satisfied, since for $\ell \neq r$ $U_{N,\ell}$ and $V_{N,r}$ commute. The relation $c_j^{\nu_j} = 1$ is satisfied by $R(c_j) = e^{2\pi i \beta_j / \nu_j}$, and the rest of the relations (7.10) are also satisfied. The remaining relation (7.12) for the orbifold braid groups implies

$$\prod_{\ell=1}^g R([a_\ell, b_\ell^{-1}]) \prod_i R(\sigma_i)^{-2} R(c_1 \cdots c_m) = 1$$

Note that from (7.16) we have $U_{N,\ell} V_{N,\ell}^{-1} = \xi_N^2 V_{N,\ell}^{-1} U_{N,\ell}$. Thus, we obtain

$$\xi_N^{2g} \cdot \xi_N^{2(n-1)} \cdot e^{2\pi i \sum_{j=1}^m \frac{\beta_j}{\nu_j}} = 1,$$

which gives $\exp(2\pi i((g+n-1)/N + \sum_j \beta_j / \nu_j)) = 1$, namely (7.17). \square

7.3. Orbifold line bundles and orbifold Euler numbers

We clarify here the relation (mentioned in Remarks 7.4 and 7.6) between the anyon representations described above and the Seifert invariants of orbifold line bundles. For a complex vector bundle \mathcal{E} of rank n over a manifold X of real dimension $2n$, the Euler number $\chi(\mathcal{E})$ is the integral on X of the Euler class $e(\mathcal{E})$. In the case of a line bundle on a 2-dimensional surface, the Euler number is the integral of the first Chern class. For an orbifold line bundle \mathcal{L} on a (good) 2-dimensional orbifold Σ , the Euler number $\chi(\mathcal{L})$ is replaced by an *orbifold Euler number* (see [62], p.437)

$$(7.18) \quad \chi^{orb}(\mathcal{L}) = \chi(\mathcal{L}) - \sum_{j=1}^m \frac{\beta_j}{\nu_j},$$

where the Euler number $\chi(\mathcal{L})$ is corrected by a contribution for each cone point x_j , $j = 1, \dots, m$ of the orbifold. These corrections are of the form β_j/ν_j , where ν_j is the order of the stabilizer $\mathbb{Z}/\nu_j\mathbb{Z}$ of the cone point x_j and the β_j are the Seifert invariants of the orbifold line bundles. These are obtained by considering the associated principal $U(1)$ -bundle P and the exact sequence

$$1 \rightarrow \mathbb{Z} \rightarrow \pi_1(P) \rightarrow \Gamma \rightarrow 1,$$

where if c_j is one of the generators of Γ with $c_j^{\nu_j} = 1$ and α is the generator of the fundamental group \mathbb{Z} of the fiber, then the β_j , with $0 \leq \beta_j \leq \nu_j - 1$, are defined by the relation

$$\alpha^{\beta_j} = \tilde{c}_j,$$

where \tilde{c}_j is a preimage of c_j in $\pi_1(P)$. By the Hopf theorem, the Euler number of a line bundle on a 2-dimensional surface is a sum over zeros of a section of the line bundle, counted with multiplicity. One can then think of the orbifold Euler number (7.18) as a modification of this counting, where additional zeros are counted at the cone points, with multiplicities β_j , but so that each zero only contributes a fraction $1/\nu_j$ of a zero at a regular point. This is consistent with the Satake orbifold Euler characteristic $\chi^{orb}(\Sigma)$, where vertices of a triangulation that are located at cone points are counted with a factor of $1/\nu_j$. In fact $\chi^{orb}(\Sigma)$ is the orbifold Euler number of the orbifold tangent bundle.

8. Laughlin-type wave functions

This section is more speculative in nature. It contains some observations on how one may naturally encounter some Laughlin-type functions in the geometric setting described in the previous sections. The Laughlin wave function can be regarded as a generalization of the Slater function

$$(8.1) \quad \Psi_{\text{Slater}}(z_1, \dots, z_n) = V(z_1, \dots, z_n) \cdot e^{-\sum_{i=1}^n \frac{|z_i|^2}{4\ell^2}},$$

with the Vandermonde determinant

$$V(z_1, \dots, z_n) = \prod_{1 \leq i < j \leq n} (z_i - z_j)$$

and with the magnetic length $\ell = \sqrt{\frac{\hbar c}{eB}}$. The Slater function describes non-interacting fermions in a magnetic field, for the full filling of the lowest Landau level. The Laughlin wave function takes the form

$$(8.2) \quad \Psi_{\text{Laughlin}}(z_1, \dots, z_n) = V(z_1, \dots, z_n)^p \cdot e^{-\sum_{i=1}^n \frac{|z_i|^2}{4\ell^2}},$$

so that it acquires a p -fold zero along the diagonals $z_i = z_j$. The exponent p is taken to be an odd integer, so that antisymmetry is preserved. In the case where p is an even integer, one considers functions of the form

$$(8.3) \quad \Psi_{\text{Pfaffian}}(z_1, \dots, z_n) = \text{Pfaff} \left(\frac{1}{z_i - z_j} \right) \cdot V(z_1, \dots, z_n)^p \cdot e^{-\sum_{i=1}^n \frac{|z_i|^2}{4\ell^2}}.$$

Explicit algorithmic methods for expressing even powers of the Vandermonde determinant as combinations of Schur functions, and Laughlin wave functions as combinations of Slater functions, are discussed in [5], [61].

By analogy with the expression $\prod_i (z_i - z)^p$ for a vortex of vorticity p centered at z , the Vandermonde determinant $V(z_1, \dots, z_n)$ in the Slater wave function can be thought of as describing cyclotron motion of n non-interacting fermions on the plane with magnetic field corresponding to a completely filled lower Landau level $\nu = 1$, and the corresponding powers $V(z_1, \dots, z_n)^p$ in the Laughlin wave function can then be thought of similarly as vortices with vorticity p , see [39], [40].

We seek here some geometric interpretation of Laughlin type wave functions related to the anyon representations described in the previous section and the geometry of the orbifold symmetric products.

8.1. Anyon representation and orbifold vector bundle

We associate orbifold vector bundles to the anyon representations described in the previous section.

Lemma 8.1. *An N -dimensional anyon representations as in Proposition 7.7 determines an orbifold line bundle \mathcal{L} on Σ , with pullback \mathcal{L}' to $\Sigma_{g'}$, and an orbifold local system V_N of (complex) rank N . These data in turn determine a rank nN orbifold vector bundle over $\text{Conf}(\Sigma, n) \subset \text{Sym}^n(\Sigma)$ of the form $\mathcal{V}_{n,N} = \mathcal{W}_n \otimes V_N$, where \mathcal{W}_n is the restriction to $\text{Conf}(\Sigma, n)$ of the orbifold vector bundle on $\text{Sym}^n(\Sigma)$ determined by the external Whitney sum $\mathcal{L}'^{\boxplus n}$ on $\Sigma_{g'}^n$.*

Proof. By construction, the anyon representation is representation of $\pi_1^{\text{orb}}(\text{Conf}(\Sigma, n))$ hence it determines an orbifold local system on $\text{Conf}(\Sigma, n)$, in the same way as representations of the ordinary fundamental group define local systems. The Seifert data β_j of the anyon representation determine an orbifold line bundle \mathcal{L} over Σ . The product $\mathcal{W}_n \otimes V_N$ is then an orbifold vector bundle over $\text{Conf}(\Sigma, n)$ of complex rank nN . \square

The orbifold Chern number of \mathcal{W}_n is obtained as follows.

Lemma 8.2. *Let \mathcal{L} be an orbifold line bundle on the good 2-dimensional orbifold $\Sigma = \Sigma_{g'}/G$. Let \mathcal{L}' be the pullback G -equivariant line bundle on $\Sigma_{g'}$. Consider the n -fold external Whitney sum $\mathcal{L}'^{\boxplus n}$ over $\Sigma_{g'}^n$, and let \mathcal{W}_n be the corresponding orbifold vector bundle over $\text{Sym}^n(\Sigma)$. Then the orbifold Euler number is given by*

$$(8.4) \quad \chi^{\text{orb}}(\mathcal{W}_n) = \frac{\chi(\mathcal{L}')^n}{n!(\#G)^n} = \frac{1}{n!} \chi^{\text{orb}}(\mathcal{L})^n.$$

Proof. By Theorem 3.6 of [62] the orbifold Euler number $\chi^{\text{orb}}(\mathcal{L})$ of an orbifold line bundle on Σ is related to the Euler number of a line bundle \mathcal{L}' on $\Sigma_{g'}$ that orbifold covers \mathcal{L} by

$$(8.5) \quad \chi(\mathcal{L}') = \chi^{\text{orb}}(\mathcal{L}) \frac{\#G}{m},$$

where m is the number of times the circle in the fiber of the principal $U(1)$ -bundle $P(\mathcal{L}')$ wraps around the circle in the fiber of $P(\mathcal{L})$. When $m = 1$, one obtains $\chi^{\text{orb}}(\mathcal{L}) = (\#G)^{-1} \chi(\mathcal{L}')$. The external Whitney sum is the Whitney sum $\oplus_{i=1}^n \pi_i^* \mathcal{L}'$, where $\pi_i : \Sigma_{g'}^n \rightarrow \Sigma_{g'}$ is the projection to the i -th factor. The

Euler class of a Whitney sum is the cup product of the Euler classes, hence we have $e(\mathcal{L}'^{\boxplus n}) = \bigwedge^n e(\mathcal{L}')$ and the Euler number is $\chi(\mathcal{L}'^{\boxplus n}) = \int_{\Sigma_{g'}^n} e(\mathcal{L}'^{\boxplus n}) = \chi(\mathcal{L}')^n$. Finally, the relation between the orbifold Euler number of \mathcal{W}_n on the symmetric products $\text{Sym}^n(\Sigma)$ and the Euler number of $\mathcal{L}'^{\boxplus n}$ on $\Sigma_{g'}^n$ is

$$\chi^{orb}(\mathcal{W}_n) = \frac{1}{\#G_n} \chi(\mathcal{L}'^{\boxtimes n}) = \frac{\chi(\mathcal{L}')^n}{n!(\#G)^n} = \frac{1}{n!} \chi^{orb}(\mathcal{L})^n.$$

□

Local systems have torsion Chern classes, hence they do not change the differential form realizing the Euler class of $\mathcal{L}'^{\boxtimes n}$ and its integration on $\Sigma_{g'}^n$.

8.2. Mathai–Quillen formalism

For a vector bundle \mathcal{E} of (real) rank $2n$ over a smooth manifold X of (real) dimension $2n$, the Euler class, whose integral $\chi(\mathcal{E}) = \int_X e(\mathcal{E})$ computes the Euler number, is the pullback along the zero section of the bundle of a representative of the Thom class. Pullbacks $e_s(\mathcal{E})$ along other sections give the same cohomology class. By Chern–Weil theory, the Euler form can be written as the Pfaffian of the curvature $\Omega^\mathcal{E}$ of a hermitian connection on the bundle \mathcal{E} ,

$$e(\mathcal{E}) = \frac{1}{(2\pi)^m} \text{Pfaffian}(\Omega^\mathcal{E}).$$

The Pfaffian of an antisymmetric matrix can be written in terms of the Berezin integral in fermionic coordinates

$$\text{Pfaffian}(A) = \int \mathcal{D}\xi \exp\left(\frac{1}{2}\xi^i A_{ij}\xi^j\right),$$

hence one can write the Euler form as

$$e(\mathcal{E}) = \frac{1}{(2\pi)^m} \int \mathcal{D}\xi \exp\left(\frac{1}{2}\xi^i \Omega_{ij}^\mathcal{E} \xi^j\right).$$

An explicit representative for the Thom class, which is exponentially decaying along the fibers (with normalized integral) and pulls back to the Euler form along the zero section is given in [52] as

$$\Phi_{MQ}(\mathcal{E}) = \frac{-\eta^2/2}{(2\pi)^m} \int \mathcal{D}\xi \exp\left(\frac{1}{2}\xi^i \Omega_{ij}^\mathcal{E} \xi^j + i\nabla\eta^i \xi_i\right),$$

where η are the fiber coordinates. It is shown in [52] that this is indeed a closed form representing the Thom class. The pullback along a nontrivial

section gives

$$(8.6) \quad e_s(\mathcal{E}) = \frac{-s^2/2}{(2\pi)^m} \int \mathcal{D}\xi \exp \left(\frac{1}{2} \xi^i \Omega_{ij}^\xi \xi^j + i \nabla s^i \xi_i \right).$$

If the section s is scaled by a factor λ , in the limit of large $\lambda \rightarrow \infty$ the form $e_s(\mathcal{E})$ localizes on the zero set $Z_s = \{x \in X \mid s(x) = 0\}$ of the section, hence recovering the Hopf theorem. The Mathai–Quillen formalism has found useful applications in physics, based on the observation [3] that the partition function of certain $N = 1$ supersymmetric gauge theories can be written as formal functional integral analog of $\int_X e_s(\mathcal{E})$.

8.3. Vandermonde determinants and symmetric products

For $S = (s_1, \dots, s_n)$, let $e_j(S)$ be the j -th elementary symmetric function in these variables, with

$$\prod_{j=1}^n (1 + ts_j) = \sum_{j=0}^n t^j e_j(S).$$

Given n symmetric polynomials f_1, \dots, f_n , the Jacobian

$$J(f_1, \dots, f_n) = \det \left(\frac{\partial f_i}{\partial s_j} \right),$$

while in the variables e_j , the Jacobian

$$J_e(f_1, \dots, f_n) = \det \left(\frac{\partial f_i}{\partial e_j} \right)$$

is related to $J(f_1, \dots, f_n)$ by

$$J(f_1, \dots, f_n) = J_e(f_1, \dots, f_n) \cdot V,$$

where

$$V(s_1, \dots, s_n) = \prod_{1 \leq i < j \leq n} (s_i - s_j) = J(e_1, \dots, e_n)$$

is the Vandermonde determinant, see [46], where explicit expressions in terms of Schur functions are given for the Jacobians J_e of complete functions and power sums. Thus, on the symmetric products $\text{Sym}^n(\Sigma)$ we should interpret the Vandermonde determinant $V(s_1, \dots, s_n)$ as the Jacobian of the change of local coordinates between the coordinates (s_1, \dots, s_n) of Σ^n to the coordinates given by the symmetric functions (e_1, \dots, e_n) .

8.4. Laughlin type functions from anyon representations

Consider an N -dimensional anyon representation as in Proposition 7.7 and the associated orbifold vector bundle $\mathcal{V}_{n,N}$ on $\text{Conf}(\Sigma, n)$, constructed as in Lemma 8.1. Let \mathcal{L} be the orbifold line bundle on Σ determined by the Seifert data of the anyon representation and \mathcal{L}' the pullback to $\Sigma_{g'}$. Let \underline{s} be a section of $\mathcal{L}'^{\boxplus n}$ determined by a n -tuple of sections of \mathcal{L}' , and let $e_{\underline{s}}(\mathcal{L}'^{\boxplus n})$ be the Mathai–Quillen representative of the Euler class. For the orbifold vector bundle \mathcal{W}_n , integration in the fiber direction now takes place with respect to coordinates given by the elementary symmetric functions in the original coordinates. This can be expressed in terms of the original coordinates by introducing the change of variables, as above, in the form of the Vandermonde determinant $V(\underline{s})$. Thus, we obtain an expression that has a product of $V(\underline{s})$ combined with an exponentially decaying factor in the fiber coordinates, as in the Slater wave function. When we further tensor with the local system V_N , each block of n -coordinates in the measure along the fiber directions acquires a factor equal to the Vandermonde determinant, while the Euler class is unchanged, hence producing a product of a power $V(\underline{s})^N$ of the Vandermonde determinant with an exponentially decaying factor in the s -coordinates, as in the Laughlin wave function. We still need to check that the power N is related to the denominators of the fractions in the Hall conductance, as is the case for the Laughlin wave functions. This is satisfied in our setting, because of the relation (7.17) in the anyon representation. In fact, notice that, for the data of an orbifold vector bundle on Σ the quantization of the Hall conductance is given, through the higher twisted index theorem, by integer multiples of Stake orbifold Euler characteristics $\chi^{orb}(\Sigma) = \chi(\Sigma_{g'})/\#G$, hence the denominator is the order of G , which is also the least common multiple of the orders ν_j of the stablizers of the cone points (Lemma 7.11 of [25]). The relation (7.17) then relates N to $\#G$.

8.5. Vandermonde determinants, Selberg integrals, and Euler characteristics

The description of the Vandermonde determinant as Jacobian of the change of coordinates on symmetric products also leads to the well known probability distributions in random matrix theory

$$(8.7) \quad \int_{\mathcal{V}_{\mathbb{R}}(n)} \exp(-\|y\|^2/2) dy = \kappa \int_{\mathbb{R}^n} V(x_1, \dots, x_n)^\alpha \exp\left(-\sum_i x_i^2/2\right) dx,$$

where $\mathcal{V}_{\mathbb{F}}(n)$ is the space of $n \times n$ -hermitian matrices, with \mathbb{F} either the real numbers, the complex numbers or the quaternions, respectively with $\alpha = \dim_{\mathbb{R}} \mathbb{F} \in \{1, 2, 4\}$. The numerical factor κ can be computed explicitly using Selberg integrals, in terms of Gamma functions, see [24] p.121.

Selberg integrals, and expectation values with respect to the random matrix probability distribution

$$V(x_1, \dots, x_n)^\alpha \exp \left(- \sum_i x_i^2 / 2 \right) dx,$$

play a crucial role in the Harer–Zagier computation of the orbifold Euler characteristic of the moduli spaces $\mathcal{M}_{g,n}$ of algebraic curves of genus g with n marked points, [36]. More recently, a *parameterized Euler characteristic* of $\mathcal{M}_{g,n}$ was introduced in [33]. This depends on a continuous parameter γ and interpolates between the case of complex and real curves, respectively corresponding to $\gamma = 1$ and $\gamma = 1/2$. The parameterized Euler characteristic is expressed in [33] in terms of Jack symmetric functions and of expectation values

$$(8.8) \quad \langle f(\lambda) \rangle = \frac{\int_{\mathbb{R}^n} f(\lambda) |V(\lambda)|^{2\gamma} e^{-\frac{\gamma}{2} p_2(\lambda)} d\lambda}{\int_{\mathbb{R}^n} |V(\lambda)|^{2\gamma} e^{-\frac{\gamma}{2} p_2(\lambda)} d\lambda},$$

with $p_2(\lambda) = \sum_{i=1}^n \lambda_i^2$. More precisely, the parameterized Euler characteristic is explicitly computed in [33] in terms of Selberg integrals of the form

$$\int_{\mathbb{R}^n} |V(\lambda)|^{2\gamma} \prod_{j=1}^n \left(1 + i \frac{\lambda_j}{a} \right)^{-\alpha} \left(1 - i \frac{\lambda_j}{b} \right)^{-\beta} d\lambda,$$

which again can be computed explicitly in terms of Gamma functions.

It would be interesting to see if Laughlin type wave functions would arise in analogous computations of orbifold Euler characteristic of moduli spaces of good 2-dimensional orbifolds. Notice that a Teichmüller theory for orbifold was developed in [68], see also [16].

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